

# Introduction to Accelerator Physics

## 2011 Mexican Particle Accelerator School

### Lecture 7/7: Some “Fun”: Corrections, Nonlinear Dynamics, Putting It All Together

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Tuesday, October 4, 2011



# Review

Hill's equation  $x'' + K(s)x = 0$

quasi – periodic ansatz solution  $x(s) = \sqrt{\epsilon\beta(s)} \cos[\phi(s) + \phi_0]$

$$\beta(s) = \beta(s + C) \quad \gamma(s) \equiv \frac{1 + \alpha(s)^2}{\beta(s)}$$

$$\alpha(s) \equiv -\frac{1}{2}\beta'(s) \quad \phi(s) = \int \frac{ds}{\beta(s)}$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_0+C} = \begin{pmatrix} \cos \Delta\phi_C + \alpha(s) \sin \Delta\phi_C & \beta(s) \sin \Delta\phi_C \\ -\gamma(s) \sin \Delta\phi_C & \cos \Delta\phi_C - \alpha(s) \sin \Delta\phi_C \end{pmatrix} \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$

betatron phase advance

$$\Delta\phi_C = \int_{s_0}^{s_0+C} \frac{ds}{\beta(s)}$$

$$\text{Tr } M = 2 \cos \Delta\phi_C$$

$$M = I \cos \Delta\phi_C + J \sin \Delta\phi_C \quad I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad J = \begin{pmatrix} \alpha(s) & \beta(s) \\ -\gamma(s) & -\alpha(s) \end{pmatrix}$$

Courant – Snyder invariant

$$\mathcal{W} \equiv A^2 = \gamma_0 x_0^2 + 2\alpha_0 x_0 x'_0 + \beta_0 x'_0{}^2$$



# General Non-Periodic Transport Matrix

- We can parameterize a general non-periodic transport matrix from  $s_1$  to  $s_2$  using the lattice parameters

$$M(s_2) = \begin{pmatrix} \sqrt{\frac{\beta(s_2)}{\beta(s_1)}} [\cos \Delta\phi + \alpha(s_1) \sin \Delta\phi] & \sqrt{\beta(s_1)\beta(s_2)} \sin \Delta\phi \\ -\frac{[\alpha(s_2) - \alpha(s_1)] \cos \Delta\phi + [1 + \alpha(s_1)\alpha(s_2)] \sin \Delta\phi}{\sqrt{\beta(s_1)\beta(s_2)}} & \sqrt{\frac{\beta(s_1)}{\beta(s_2)}} [\cos \Delta\phi - \alpha(s_2) \sin \Delta\phi] \end{pmatrix}$$

- This does not have a pretty form like the periodic matrix

However both can be expressed as

$$M = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix}$$

where the C and S terms are cosine-like and sine-like; the second row is the s-derivative of the first row!

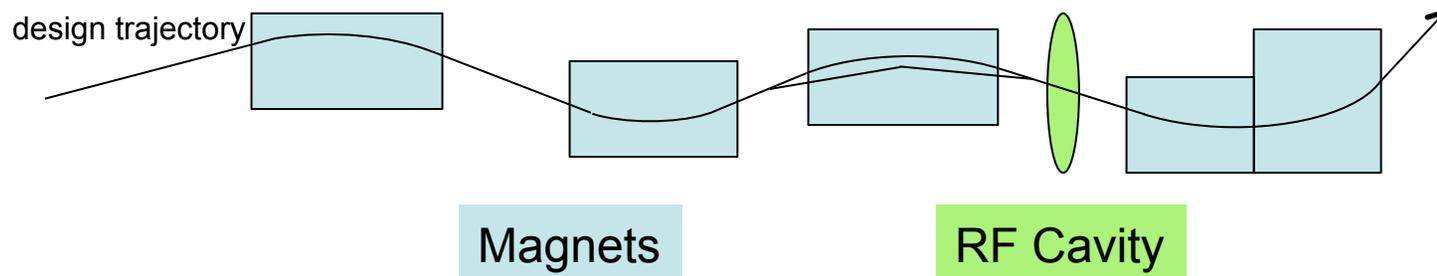
The most common use of this matrix is the  $m_{12}$  term:

$$\Delta x(s_2) = \sqrt{\beta(s_1)\beta(s_2)} \sin(\Delta\phi) x'(s_1)$$

Effect of angle kick  
on downstream position



# Design Orbit Perturbations



- Sometimes need a local change  $\Delta x(s)$  to the design orbit
  - But we really only get changes in angle  $\Delta x'$  from magnets
  - e.g. small dipole “corrector”:  $\Delta x' = B_{\text{corrector}} L_{\text{corrector}} / (B\rho)$
  - Changes to/corrections of design orbit from dipole correctors
  - Linear errors add up via linear superposition

$$\begin{pmatrix} \Delta x(s_2) \\ \Delta x(s_1) \end{pmatrix} = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix} \begin{pmatrix} 0 \\ \Delta x'(s_1) \end{pmatrix}$$

$$\Delta x(s_2) = \Delta x'(s_1) \sqrt{\beta(s_1)\beta(s_2)} \sin \Delta\phi_{12}$$

$$\Delta x'(s_2) = \Delta x'(s_1) \sqrt{\frac{\beta(s_1)}{\beta(s_2)}} [\cos \Delta\phi_{12} - \alpha(s_2) \sin \Delta\phi_{12}]$$



## Two-Bump

$$\Delta x(s_2) = \Delta x'(s_1) \sqrt{\beta(s_1)\beta(s_2)} \sin \Delta\phi_{12}$$

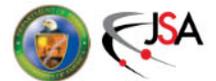
$$\Delta x'(s_2) = \Delta x'(s_1) \sqrt{\frac{\beta(s_1)}{\beta(s_2)}} [\cos \Delta\phi_{12} - \alpha(s_2) \sin \Delta\phi_{12}]$$



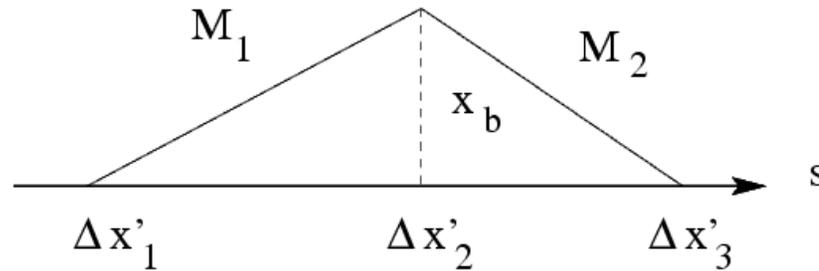
- But this orbit error now changes all later positions and angles
  - Add another dipole corrector at a location where  $\Delta\phi_{12} = k\pi$   
At this point the distortion from the original dipole corrector is all  $x'$  that we can cancel with the second dipole corrector.

$$\Delta x'(s_2) = \Delta x'(s_1) \sqrt{\frac{\beta(s_1)}{\beta(s_2)}} + \text{angle from } s_2 \text{ dipole}$$

- Called a **two-bump**: localized orbit distortion from two correctors
- But requires  $\Delta\phi_{12} = k\pi$  between correctors



# Three-Bump



- A general local orbit distortion from three dipole correctors
  - Constraint is that net orbit change from sum of all three kicks must be zero

$$\begin{pmatrix} C_2 & S_2 \\ C'_2 & S'_2 \end{pmatrix} \left[ \begin{pmatrix} C_1 & S_1 \\ C'_1 & S'_1 \end{pmatrix} \begin{pmatrix} 0 \\ \Delta x'_1 \end{pmatrix} + \begin{pmatrix} 0 \\ \Delta x'_2 \end{pmatrix} \right] + \begin{pmatrix} 0 \\ \Delta x'_3 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\Delta x'_1 = \frac{x_b}{S_1} \quad \Delta x'_2 = -\frac{C_2 S_1 + S_2 S'_1}{S_1 S_2} x_b \quad \Delta x'_3 = \frac{S_2}{S_1^2} x_b$$

- Bump amplitude  $x_b = S_1 \Delta x'_1$
- Only **three-bump** requirement is that  $S_1, S_2 \neq 0$



# Steering Error in Synchrotron Ring

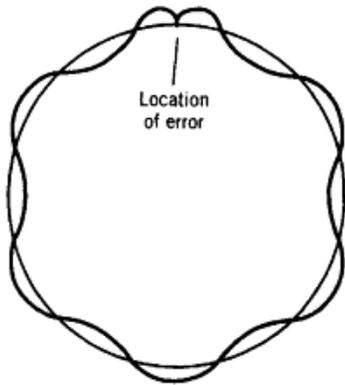
- Short steering error  $\Delta x'$  in a ring with periodic matrix  $M$ 
  - Solve for new periodic solution or design orbit  $(x_0, x'_0)$

$$M \begin{pmatrix} x_0 \\ x'_0 \end{pmatrix} + \begin{pmatrix} 0 \\ \Delta x' \end{pmatrix} = \begin{pmatrix} x_0 \\ x'_0 \end{pmatrix}$$

- Note that  $(x_0=0, x'_0=0)$  is not the periodic solution any more!

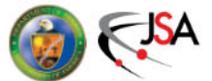
$$\begin{pmatrix} x_0 \\ x'_0 \end{pmatrix} = (I - M)^{-1} \begin{pmatrix} 0 \\ \Delta x'_0 \end{pmatrix}$$

$$\begin{aligned} (I - M)^{-1} &= (I - e^{(2\pi Q)J})^{-1} = ([e^{\pi Q J} (e^{-\pi Q J} - e^{\pi Q J})])^{-1} \\ &= -(2J \sin(\pi Q))^{-1} (e^{\pi Q J})^{-1} \\ &= \frac{1}{2 \sin(\pi Q)} (J \cos(\pi Q) + I \sin(\pi Q)) \end{aligned}$$



$$x_0 = \frac{\Delta x'_0}{2} \tan(\pi Q) \quad \rightarrow \infty \text{ if } Q = k\pi$$

integer resonances



# Focusing Error in Synchrotron Ring

- Short focusing error in a ring with periodic matrix  $M$ 
  - Now solve for  $\text{Tr } M$  to find effects on tune  $Q$

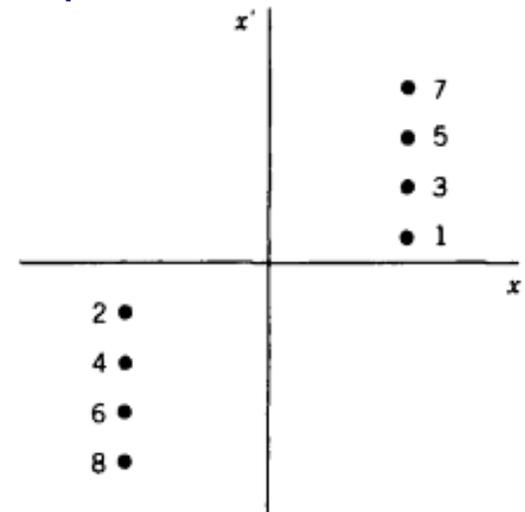
$$M_{\text{new}} = M \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix}$$

$$\frac{1}{2} \text{Tr } M = \cos(2\pi Q_{\text{new}}) = \cos(2\pi Q_0) - \frac{1}{2} \frac{\beta_0}{f} \sin(2\pi Q_0)$$

- For small errors  $Q_{\text{new}} = Q_0 + \Delta Q$  we can expand to find

$$\Delta Q \approx \frac{1}{4\pi} \frac{\beta_0}{f}$$

- Quadrupole errors also cause resonances when  $Q = k/2$ : **half-integer resonances**



# Chromaticity Correction

Natural chromaticity  $\xi_N \equiv \left( \frac{\Delta Q}{Q} \right) / \left( \frac{\Delta p}{p_0} \right) = -\frac{1}{4\pi Q} \oint K(s)\beta(s) ds$

- How can we control chromaticity in our synchrotron ring?
  - We need a way to connect momentum offset  $\delta$  to focusing
  - Dispersion (momentum-dependent position) and sextupoles (nonlinear focusing depending on position) come to rescue

$$x(s) = x_{\text{betatron}}(s) + \eta_x(s)\delta$$

Sextupole B field  $B_y = b_2 x^2$

$$B_y(\text{sext}) = b_2 [x_{\text{betatron}}(s) + \eta_x(s)\delta]^2 \approx b_2 x_{\text{betatron}}^2 + 2b_2 x_{\text{betatron}}(s)\eta_x(s)\delta$$

Nonlinear!                      like a quadrupole K(s)!

- Total chromaticity from all sources is then

$$\xi = -\frac{1}{4\pi Q} \oint [K(s) - b_2(s)\eta_x(s)] ds$$

- Strong focusing (large K) requires large sextupoles, **nonlinearity!**

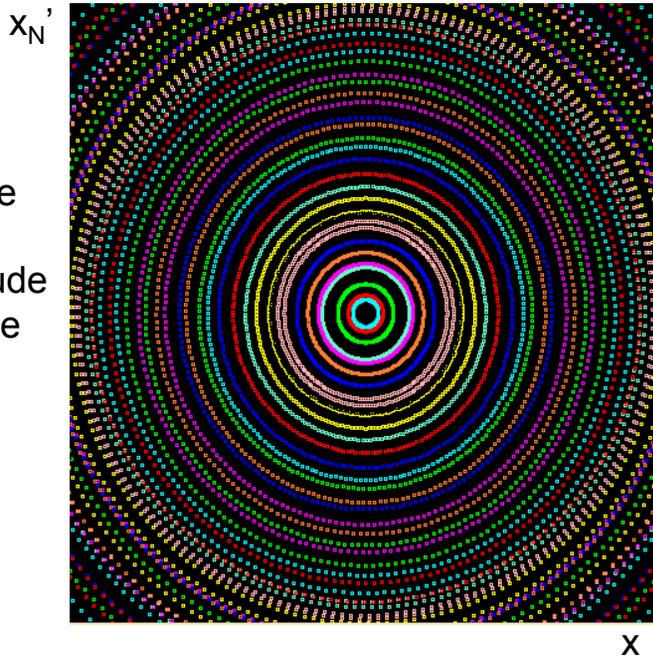


# How Bad Could It Be?

- We expanded so far assuming nonlinearities are small
- How bad could it really be?
  - Even **one** large nonlinear (sextupole, octupole, ...) error in a ring can drive **all** nonlinear resonances  $kQ_x + lQ_y = m$
  - Large-amplitude particle motion will almost invariably be **unstable**, even **chaotic** in the true dynamical sense

Q=10.205  
Linear  
No sextupole

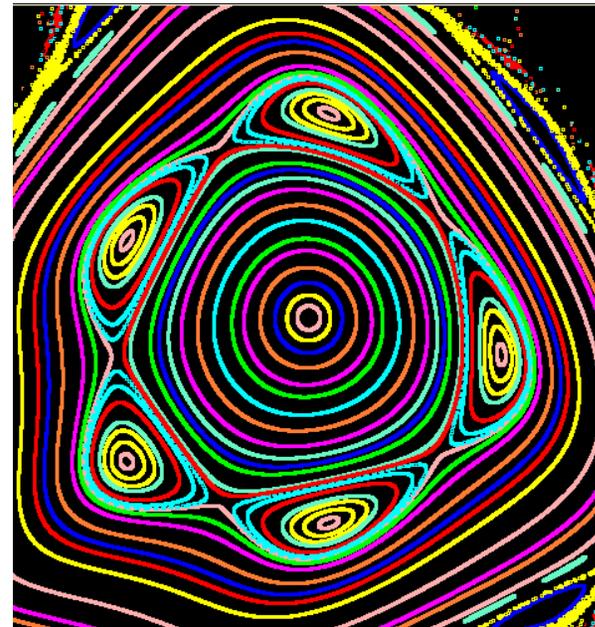
High amplitude  
motion stable



Q=10.205  
**Nonlinear**  
**One** sextupole  
 $b_2=0.564$

High amplitude  
motion unstable

$5Q_x=m$   
resonance  
islands

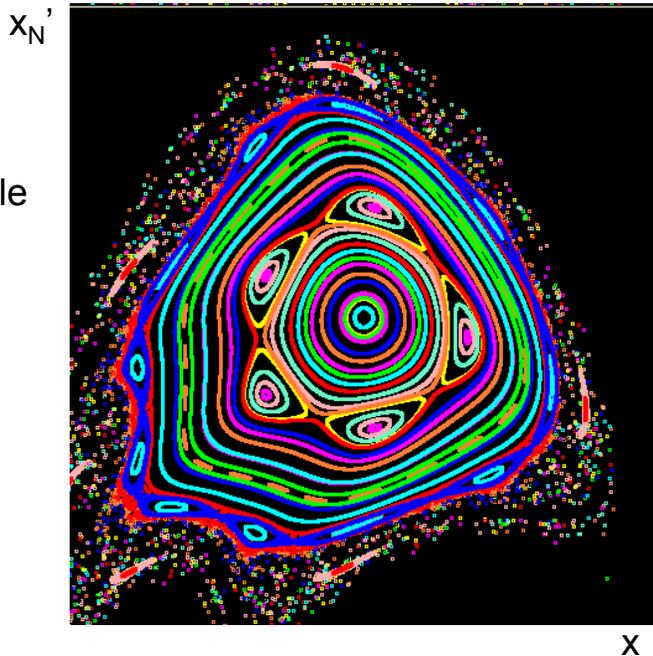


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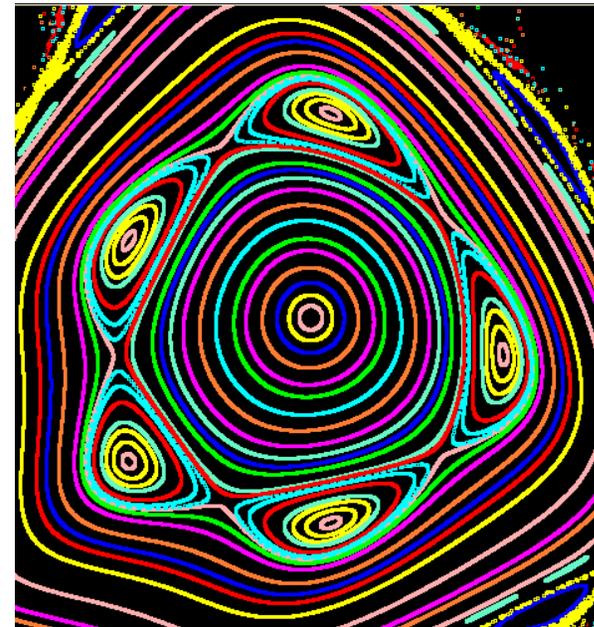
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Q=10.205  
**Nonlinear**  
 One sextupole  
 $b_2=1041$



Q=10.205  
**Nonlinear**  
 One sextupole  
 $b_2=0.564$



High amplitude  
 motion unstable

$5Q_x=m$   
 resonance  
 islands



# Tune Diagram

- We often plot  $(Q_x, Q_y)$  on a tune diagram  
Also plot resonance lines  $kQ_x + lQ_y = m$

We can still analyze nonlinear resonances with perturbation theory

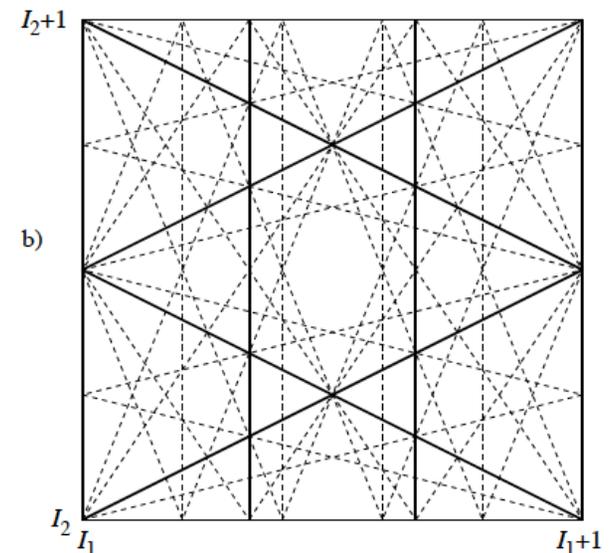
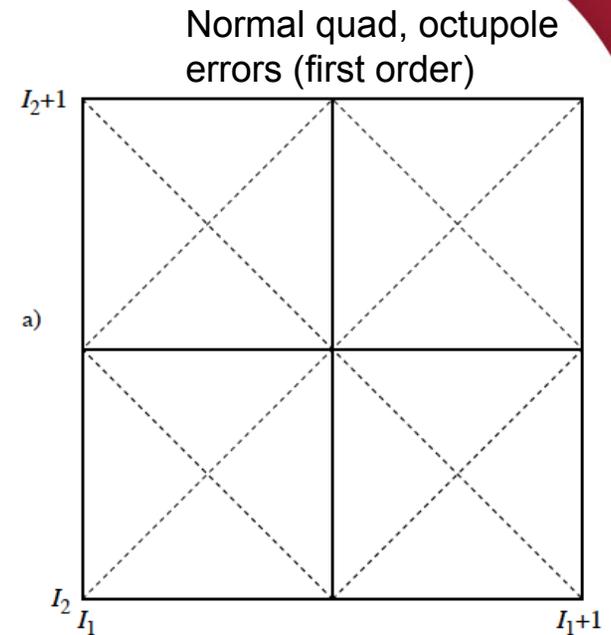
Various resonances are driven to various “orders” in perturbative expansions

If nonlinearities are kept small, then higher order nonlinearities can be neglected

**Thanks for the great magnets, Cherryl and Liz!**

Remember tune here is a blob, not a point!

Beam tunes  $(Q_x, Q_y)$  have distributions



Normal sextupole, decapole errors (first order)



# Hamiltonians

- The dynamical treatment of nonlinearities particle accelerators is a broad subfield of its own
  - Predictions of nonlinear behavior directly affect specifications, cost, and even feasibility
  - Cost of facility scales with magnet aperture, dynamic aperture
  - Interactions with other fields (e.g. astrophysics, Laskar)
- Perturbative approach is usually done using Hamiltonians
  - Matrices are not very convenient for analysis of nonlinearities

$$H(x, x'; s) = \frac{x'^2}{2} + K(s) \frac{x^2}{2} \quad \text{Quadratic quadrupole focusing!}$$

Hamilton's Equations

$$x' = \frac{\partial H}{\partial x'} \quad (x')' = -\frac{\partial H}{\partial x}$$

⇒ Hill's Equation again

$$x''(s) + K(s)x(s) = 0$$



# Action-Angle Coordinates

- Certain types of coordinate transforms keep Hamilton's Equations for the dynamics
  - These are called **canonical transformations** in dynamics
  - Can be generated from **generating functions**
- One very useful coordinate transformation is to **action-angle** coordinates
  - We want to treat the betatron phase  $\phi$  as a coordinate
  - Its canonical conjugate variable is an action  $J$

$$J = \frac{x^2}{2\beta(s)} \sec^2 \phi(s) = \frac{1}{2\beta(s)} [x(s)^2 + (\beta(s)x'(s) + \alpha(s)x(s))^2]$$

$$\begin{pmatrix} x \\ x' \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta(s)} \cos \phi(s) \\ -\sqrt{1/\beta(s)} [\sin \phi(s) + \alpha(s) \cos \phi(s)] \end{pmatrix}$$

- This looks horrible, but now the Hamiltonian is simpler!



# Action-Angle Coordinates

- How simple is the action-angle Hamiltonian?

$$H(J, \phi; s) = \frac{J}{\beta(s)}$$

- Periodic in  $s$  (like Hill's Equation)
- No dependence on  $\phi$  so  $J$  is a constant of the motion
- Integrating this over one turn gives a “one-turn” Hamiltonian

$$H_{\text{one turn}} = J \oint \frac{ds}{\beta(s)} = 2\pi Q J$$

- Hamilton's equations give  $\Delta\phi_{\text{one turn}} = 2\pi Q$   $\Delta J_{\text{one turn}} = 0$
- The real power is that we can add nonlinear perturbation potentials to the Hamiltonian for nonlinear fields and solve in exactly the same ways

$$H(x, x', y, y'; s) = \frac{x'^2}{2} + K_x(s) \frac{x^2}{2} + \frac{y'^2}{2} + K_y(s) \frac{y^2}{2} + V(x, y; s)$$

$$H(J_x, \phi_x, J_y, \phi_y; s) = \frac{J_x}{\beta_x} + \frac{J_y}{\beta_y} + V(J_x, \phi_x, J_y, \phi_y; s)$$



## “Simple” Nonlinearity: A Single Sextupole

$$H(x, x', y, y'; s) = \frac{x'^2}{2} + K_x(s) \frac{x^2}{2} + \frac{y'^2}{2} + K_y(s) \frac{y^2}{2} + \frac{b_2}{6} (x^3 - 3xy^2)$$

one dimensional ( $y = 0$ )  $H(x, x'; s) = \frac{x'^2}{2} + K_x(s) \frac{x^2}{2} + \frac{b_2}{6} x^3$

- The extra term will add dependencies on  $\phi$  so J is no longer a constant of the motion: **phase space distortion**
- Using  $x = \sqrt{2J_x\beta_x} \cos \phi_x$  we can expand the nonlinear term

$$\text{nonlinear term } V = \frac{b_2}{6} x^3 = \frac{b_2}{6} (2J_x\beta_x)^{3/2} \cos^3 \phi_x$$

$$\cos^3 \phi = \frac{1}{4} (3 \cos \phi + \cos 3\phi)$$

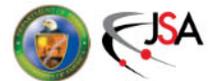
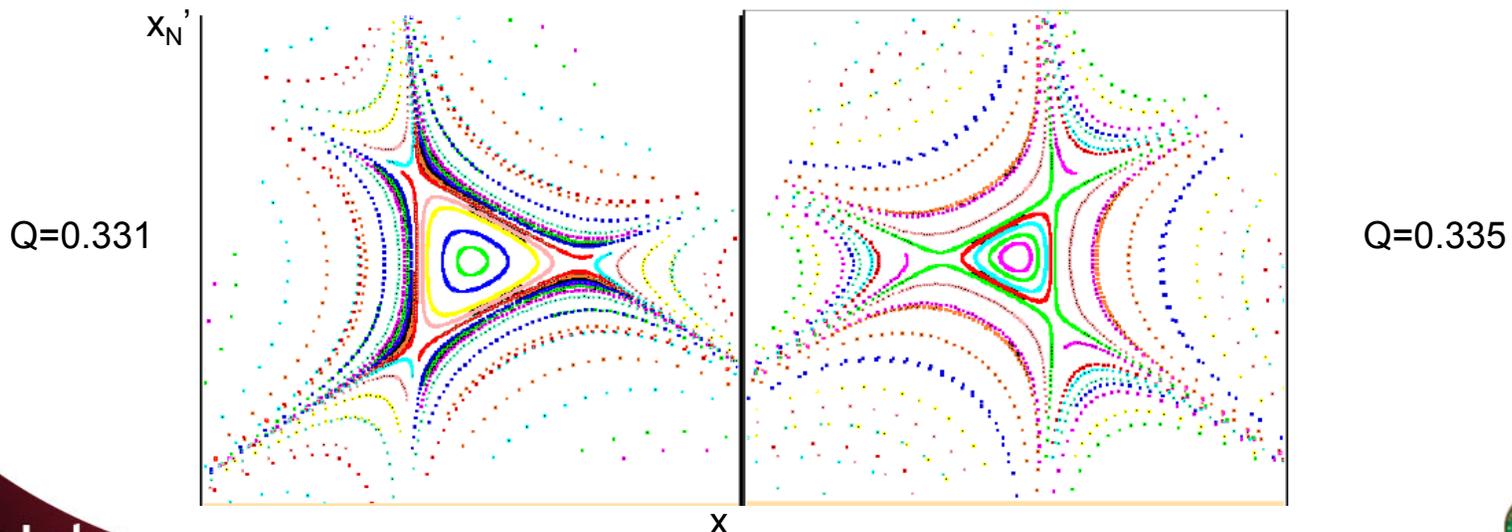
$$\text{nonlinear term } V = \frac{b_2}{6} x^3 = \frac{b_2\sqrt{2}}{12} (J_x\beta_x)^{3/2} (3 \cos \phi_x + \cos 3\phi_x)$$

- Betatron phase dependences drive  $Q_x=k$ ,  $3Q_x=k$  resonances



# Sextupole $3Q_x=k$ Resonance

- For tunes close to  $3Q_x=k$ , we can solve for the motion
  - Often just iterate Hamilton's equations to plot them
  - There end up being three fixed points of this "map" at three betatron phases separated by  $2\pi/3$  (120 degrees)
  - These three fixed points also end up being locally unstable
    - Nearby motion is hyperbolic rather than elliptical
  - Area of stable particle motion is distorted and reduced
    - So don't operate near  $3Q_x=k$  with strong sextupoles!



# Putting It All Together

- Some thoughts from 20 years in accelerators and seeing several new accelerator projects start up
- Organize
  - Users, physicists, engineers, funding groups/agencies
- Discuss
  - What sort of facility is best for needs? Initial specifications
- Write
  - Iterate more and more details of facility design, engineering
  - Eventually becomes a “conceptual design” report
  - Balance risk, feasibility, cost, benefits
- Sell
  - Pitch the project to funding agencies



# Organize

- You have an excellent start towards organizing a team
- Users are your customers!
  - Their needs must be met to politically support your project
  - Some must be “talked back down to Earth” – or talked up!
  - Identifying science benefits for facility organizes priorities
  - They primarily provide the **goals** of the facility
- Scientists and engineers
  - You have an excellent start here, particularly many students!
  - Have them travel and learn from other experts, facilities
  - Interact with users to build a strong community, communication
  - Primarily provide the **technical constraints** of the facility
- Funding agencies
  - **Sometimes** good to have them involved from early stages
  - They can provide important early **resource constraints**



# Discuss

- Iterate ideas among all groups
  - Negotiations provide a framework for communication, trust
  - Developing teamwork
  - Writing justification document for facility need
- Small group teams naturally develop
  - Focus on one subset of problem or particular area of challenge, technical risk
  - Common areas of expertise and interest
  - Teams remain in place to contribute to conceptual design
- Strong central management
  - Integrates input from teams, guides overall vision
  - Stays focused on end goal, but flexible in face of change
  - Makes critical decisions to establish baseline parameters



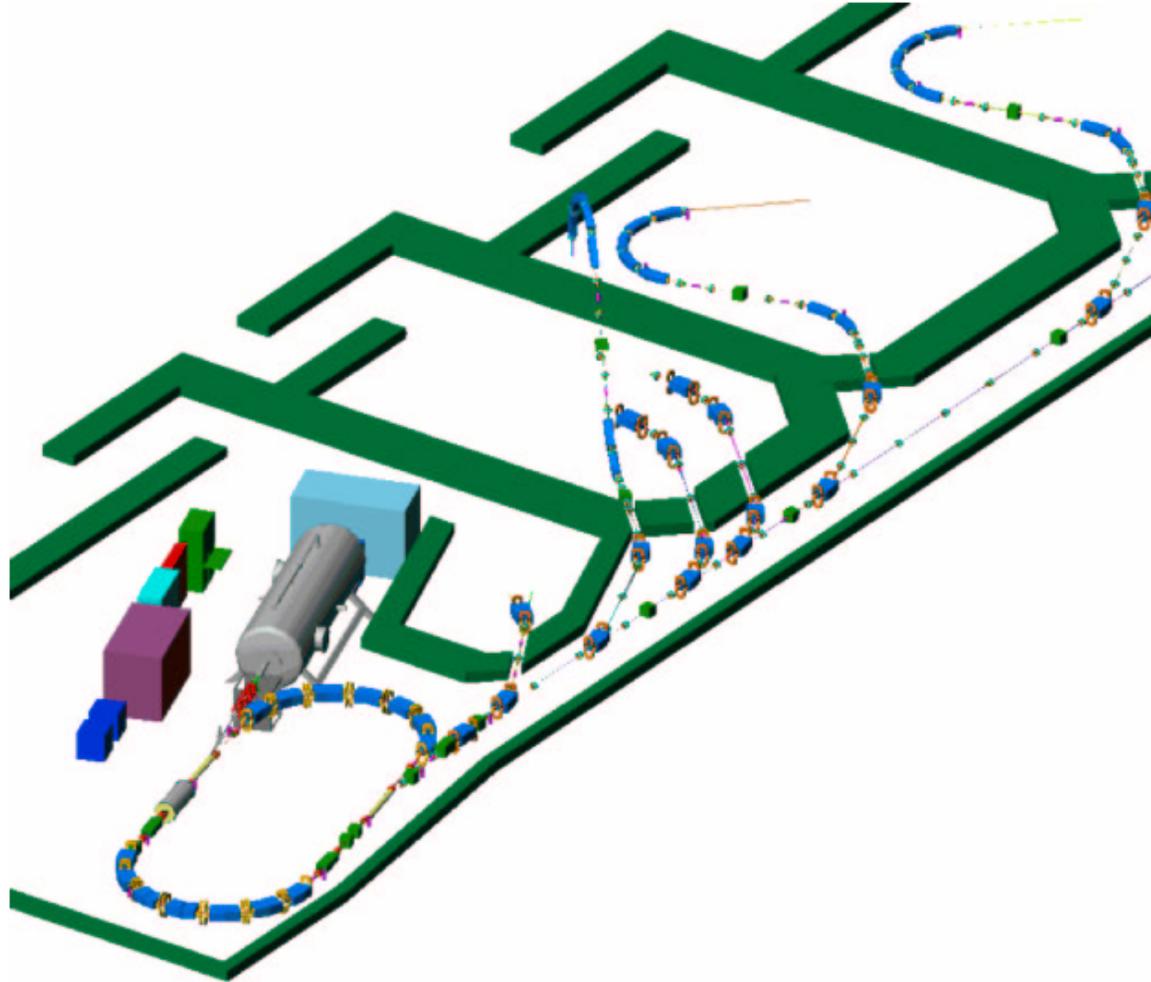
# Write

- Conceptual Design Report (CDR)
  - An assembly of physics and engineering specifications of the facility
  - Not necessarily “build to print” or “build to spec” yet
  - No detailed costing, site specifics
- The CDR should clearly identify
  - Physics goals and user community of the facility
  - All parameters necessary to meet those goals
  - How those parameters are technically produced
    - Identify areas of risk and/or required R&D
    - Identify areas of negligible risk and off-shelf availability
  - Breadth of subsystem physics/engineering design
    - Magnets, lattice, RF, injector, vacuum, cryogenics, ...
    - Desired beam properties, brilliance, number of beamlines...



# Medical Accelerator CDR Cover

## Conceptual Design of the RCMS



# Pitch

- The CDR should convince other people “we know how to build this”
  - Including identifying risks and benefits of R&D
  - CDR serves as a central focus of team efforts to remain consistent across the project: a good management tool
- CDR can also convince others “we know why we want to build this”
  - Sometimes a physics/benefit justification section is included
  - But this is often a separate document used to justify and mobilize the resources necessary to produce the CDR

