

# #9: SEXTUPOLES & CHROMATICITY

1/28/19

CATCH-22: Sextupoles are (usually) necessary to avoid NON-linear resonances that are often driven by sextupoles themselves!

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Recall motion of a particle with  $\delta = \frac{\Delta p}{p}$  (constant)

Hill's Equations

$$\begin{cases} x'' + \frac{k}{(1+\delta)} x = 0 \\ y'' - \frac{k}{(1+\delta)} y = 0 \end{cases} \left\{ \begin{array}{l} k \text{ "weaken"} \\ \text{with } \delta \end{array} \right.$$

- CHROMATICITY measures the rate of change of tunes  $Q_x$  &  $Q_y$  with momentum

$$\chi \equiv \frac{dQ}{d\delta}$$

$$M \text{ e } V$$

EG 1

FODO LATTICE

$N$  identical cells, each with  $\phi$  phase advance

$$Q = \frac{N\phi}{2\pi}$$

so the NATURAL (uncorrected) chromaticity is

$$\chi_{\text{NAT}} = \frac{N}{2\pi} \frac{d\phi}{d\delta}$$

but following Eq. 3.48

$$\frac{d\phi}{d\delta} = -2 \tan(\phi/2)$$

so

$$\chi_{x, \text{NAT}} = \chi_{y, \text{NAT}} = -Q \frac{\tan(\phi/2)}{\phi/2} \sim -Q$$

This RULE OF THUMB often holds at factors of 2 level

EG 2 RHIC

$$\chi_{\text{NAT}} \approx -50$$

RMS momentum spread  $\frac{\sigma_p}{p} = \langle \delta^2 \rangle^{\frac{1}{2}} \approx 2 \times 10^{-3}$

UNCORRECTED  
tune spread

$$\sigma_Q = |\chi_{\text{NAT}}| \frac{\sigma_p}{p} = 0.1 \quad !!$$

UNACCEPTABLY shrinks available locations  
in the tune plane

$\Rightarrow$  MUST DECREASE  $|\chi|$

EG3 : SEXTUPOLE next to EVERY QUADRUPOLE

FIRST separate closed orbit + betatron oscillations:

$$X_{TOT} = \eta \delta + x$$

THEN expand the composite quad-sext kick:

$$\Delta x'_{TOT} = -q(1-\delta)(\eta\delta + x) - S(1-\delta)(\eta\delta + x)^2 + \dots$$

AND collect terms in  $x^n \delta^m$

$$\begin{aligned} (\Delta \eta') \delta + \Delta x' &= -q \cdot \eta \delta \\ &\quad - q x - S x^2 \\ &\quad + q \cdot x \delta - S 2\eta x \delta + S x^2 \cdot \delta \\ &\quad + O(\delta^2) \dots \end{aligned}$$

$x^0 \delta^1$   
 $x^m \delta^0$  GEOMETRIC

TO ORDER  $x' \delta^m$

$$\Delta x' = -q x + (q - 2\mu S) x \delta$$

So the effective "quadrupole" strength is independent of  $\delta$  IF

$$S = \frac{q}{2\mu}$$

at every quad-sext pair

- In a FODO lattice  $\mu_F \approx 2\mu_D$  so

$$|S_F| \approx \frac{1}{2} |S_D|$$

Q: What happens if  $\mu \neq 0$ ?

[Q: What do these geometric terms do?] later ...

# EG 4: CHROMATICITY CORRECTION: GENERAL CASE

Previously Equ. 8.29: If a quad increases strength by  $\Delta q$

$$\Delta Q = \frac{\beta_c}{4\pi} \cdot \Delta q$$

QUAD STRENGTH weakens with  $s$

$$\frac{dq}{ds} = -q$$

So natural chromaticities:

$$\begin{pmatrix} \chi_x \\ \chi_y \end{pmatrix}_{\text{NAT}} = -\frac{1}{4\pi} \begin{pmatrix} \sum q \beta_x \\ \sum -q \beta_y \end{pmatrix}$$

(Thick quads:  
use integrals)

Where sums  $\sum$  are over all (thin) quads.

Strength  $q$  is usually +ve (-ve) when  $\beta_x$  ( $\beta_y$ ) is large

$\Rightarrow$  NATURAL CHROMATICITIES ARE NEGATIVE

SEXTUPLES contribute a QUTD component

$$\frac{dq}{d\delta} = 2\eta S$$

So ONE FAMILY of sextupoles changes chromaticities

$$\Delta \begin{pmatrix} \chi_x \\ \chi_y \end{pmatrix} = \frac{S_i}{2\pi} \begin{pmatrix} \sum \eta \beta_x \\ \sum -\eta \beta_y \end{pmatrix} \quad \begin{cases} \mu_x \text{ is H} \\ \mu_y = 0 \end{cases}$$

in different directions

TWO FAMILIES near F + D are quadrupoles

$$\begin{pmatrix} \chi_x \\ \chi_y \end{pmatrix}_{\text{GOAL}} = 2 \begin{pmatrix} \sum_F \eta \beta_x & \sum_D \eta \beta_x \\ \sum_F -\eta \beta_y & \sum_D -\eta \beta_y \end{pmatrix} \begin{pmatrix} S_F \\ S_D \end{pmatrix} + \begin{pmatrix} \chi_x \\ \chi_y \end{pmatrix}_{\text{NAT}}$$

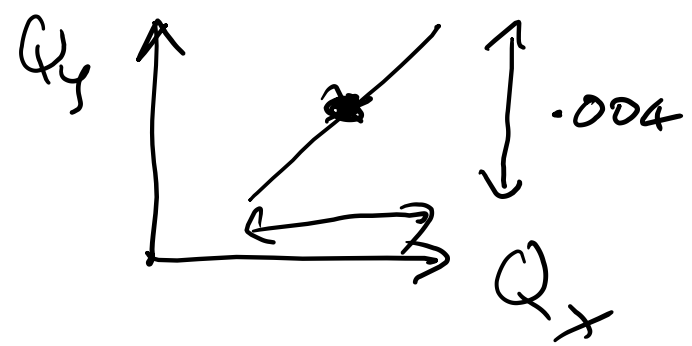
SOLVE for  $S_F$  +  $S_D$  by inverting the matrix !!

EGS RHIC (again)

$$\chi_{\text{GOAL}} = +2$$

$$\frac{\sigma_P}{P} = 2 \times 10^{-3}$$

$$\Rightarrow |\chi_{\text{GOAL}}| \frac{\sigma_P}{P} \Rightarrow \sigma_Q = .004$$



Q: Why is TUNE FOOTPRINT a blob not a line?

Q: Why are  $\chi$  goals slightly positive?  
(collective effects)

[Q: What drives the resonance lines?]



# HENON MAP : GEOMETRIC TERMS ( $\delta=0$ )

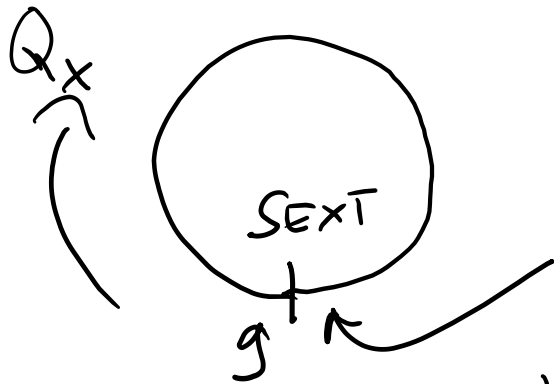
"... exhibits ALL TYPICAL PROPERTIES of... dynamical systems"  
until finished {

$$\begin{pmatrix} z \\ z' \end{pmatrix} = \begin{pmatrix} \cos(2\pi Q) & \sin(2\pi Q) \\ -\sin( ) & \cos( ) \end{pmatrix} \begin{pmatrix} z \\ z' \end{pmatrix} + \begin{pmatrix} 0 \\ z^2 \end{pmatrix}$$

The ONLY control parameter  $Q$  is like the (horizontal) tune

The nonlinear kick  $z^2$  is like a sextupole.

# ONE SEXTUPOLE (horizontal motion)



Reference point is JUST BEFORE the sextupole

KICK:  $\begin{pmatrix} x \\ x' \end{pmatrix}_{0+c} = \begin{pmatrix} x \\ x' - gx^2 \end{pmatrix}_0$

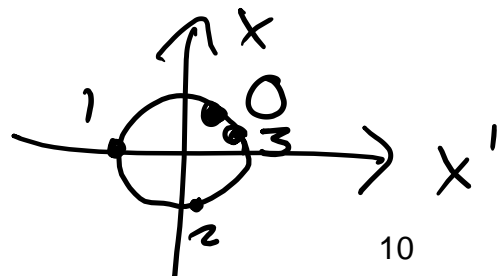
ROTATE:  $\begin{pmatrix} x \\ x' \end{pmatrix}_1 = R(2\pi Q) \begin{pmatrix} x \\ x' \end{pmatrix}_{0+c}$

EG 1

TUNE IS CLOSE TO  $1/3$

$$Q = \frac{1}{3} + \delta Q$$

NET MOTION after 3 turns  
is SMALL . . . .



$$\begin{pmatrix} x \\ x' \end{pmatrix}_3 - \begin{pmatrix} x \\ x' \end{pmatrix}_0 \approx 3\mu \begin{pmatrix} x \\ -x' \end{pmatrix}_0 - g \left[ \begin{pmatrix} s_3 \\ c_3 \end{pmatrix} x_0^2 + \begin{pmatrix} s_2 \\ c_2 \end{pmatrix} x_0^2 + \begin{pmatrix} s_1 \\ c_1 \end{pmatrix} x_0^2 \right]$$

where  $\mu = 2\pi \delta Q$ ,  $c_k = \cos\left(k \frac{2\pi}{3}\right)$ ,  $s_k = \sin\left(k \frac{2\pi}{3}\right)$ .

MORE SUCCINCTLY: 3-turn discrete

KOBAYASHI HAMILTONIAN:

$$H_3 = \frac{\mu}{2} (x^2 + x'^2) + \frac{g}{9} \sum_{k=1}^3 (c_k x + s_k x')^3$$

$$\Delta x \approx \frac{\partial H_3}{\partial x'} \cdot \Delta t$$

$$\Delta x' \approx - \frac{\partial H_3}{\partial x} \cdot \Delta t$$

with  $\Delta t = 3$

[Q: Cross terms?]

-  $H_3$  is approximately constant along a trajectory

- Factorising  $H_3$  gives

$$H_3 = \frac{g}{3} \left( \frac{2M}{g} \right)^3 + \frac{g}{12} \left( x + \sqrt{3} x' + \frac{4M}{g} \right) \left( x - \sqrt{3} x' + \frac{4M}{g} \right) \left( x - \frac{2M}{g} \right)$$

So, if  $H_3 = \frac{g}{3} \left( \frac{2M}{g} \right)^3$

then

3 straight lines !!

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EG  $Q = 0.324$

- 3 straight lines!
- Regular "circular" motion near origin
- Rapid escape along divergent arms
- 3 outer "islands" unpredicted by perturbation theory

EG  $Q = .2516 \approx \frac{1}{4}$

- 4 RESONANCE ISLANDS!
- Island hopping: only some phases are accessible.
- Chains of many small islands also appear

$Q = \frac{P}{2\pi} \approx 33?$  (large number).

EQ  $Q = 0.211 \approx 1/5$

- CHAOTIC "fly specks" DOTS are visually distinctive

- 5 islands are more rotationally symmetric

EQ  $Q = 0.185 \approx 1/6$

6 smaller islands, more rotationally symmetric,  
closer to the DYNAMIC APERTURE

# HENON'S TAXONOMY of typical behavior:

## ① REGULAR NON-RESONANT

- Roughly circular motion
- Enough dots (turns) make "continuous" lines

## ② REGULAR RESONANT

- Island hopping: only some phases accessible

## ③ RAPIDLY DIVERGENT ( $Q = -32f$ )

- Amplitude increases rapidly & regularly

## ④ CHAOS

- Visually distinctive, but WHAT IS IT?

# LYAPUNOV EXPONENT

- Start 2 particles VERY close together  
(e.g. one digit apart in double precision)

Q: How do they BEGIN to diverge in time?

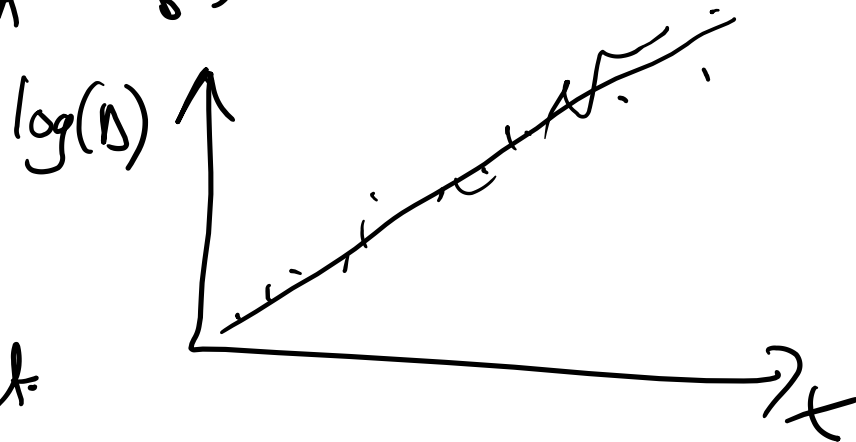
A1: LINEARLY  $\Rightarrow$  Regular

A2: EXPONENTIALLY  $\Rightarrow$  Chaotic

e.g.  $D = (x_A - x_B)^2 + (x'_A - x'_B)^2$

$$D(t) \sim D_0 e^{\lambda t}$$

Lyapunov exponent





# DYNAMIC APERTURE (DA)

- IN SIMULATION there is (usually) a clear cut maximum amplitude beyond which particles rapidly get lost - hit the beam pipe.
- IN REALITY there is noise/excitation/diffusion drives particles across the DA
- beam current decreases with time

EVEN IN 1-D (not 3-D) WITH 1 SEXTUPOLE  
this is a complex story (and NATIVE)