

# USPAS Accelerator Physics 2019 (Virtually) Texas A&M University

## Chapter 5: Action and Emittance

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# 01/29/21 L6 Ch5 Emittances and Phase Space

“One particle, many, or none?”

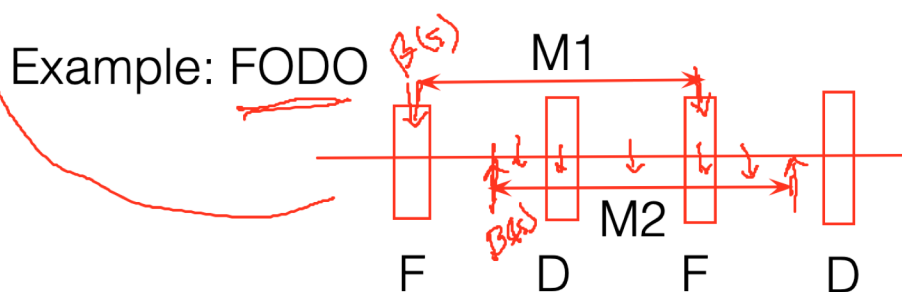
Recall the parameterization of PERIODIC linear motion (Tue slide 11)

$$M = \begin{pmatrix} \cos \mu + \alpha(s) \sin \mu & \beta(s) \sin \mu \\ -\gamma(s) \sin \mu & \cos \mu - \alpha(s) \sin \mu \end{pmatrix} \quad \mu = 2\pi Q$$

$x$  or  $y$

det M = 1       $\alpha(s) \equiv -\frac{1}{2} \frac{d\beta(s)}{ds}$        $\gamma(s) \equiv \frac{1 + \alpha^2(s)}{\beta(s)}$

Tr M = 2 cos  $\mu$        $\mu$ : does NOT depend on s



$M_1$  &  $M_2$  are both periodic of the above form

$M_1$  has  $\beta(s), \alpha(s), \gamma(s)$  at center of F

$M_2$  has " " " between F, D

Both have the same  $\mu$

2

1 particle, many particles, or no particles?

No particles: Optics (Twiss parameters, or C-S parameters) are defined by the lattice (magnets)

optical properties of lattice (periodic so far)

BUT

What about beamlines? Linacs?

Non-periodic systems?? Are there useful parameterizations of linear motion:  $\beta, \alpha, \gamma$ ?

A: yes (sorta)  $\Rightarrow$  here people usually speak more of the beam

STILL:  $\begin{pmatrix} x \\ x' \end{pmatrix}_{\text{OUT}} = M \begin{pmatrix} x \\ x' \end{pmatrix}_{\text{IN}}$   $\odot$  parameterize  $(x, x')$  ...

### 3 Single-particle coordinates in action-angle coordinates

Hill's EQN: (1D)  $x'' + K(s)x = 0 \dots$   
 suggests pseudo-SHO solutions

Ansatz: (A)  $x_n = \sqrt{2J_x \beta_x} \sin(\phi_x)$   $J_x = \text{action}$   
 $x'_n = \sqrt{\frac{2J_x}{\beta_x}} [\cos(\phi_x) - \alpha_x \sin(\phi_x)]$   $\phi = \text{angle}$  Hamiltonian

For  $n$ -turn parameterization  $J_x = \text{constant}$   
 "phase advance"  $\Leftrightarrow \phi_x$  advances

For ring  $w$ , tune  $Q_x$ ,  $\phi_{x,n} = \phi_{0x} + \underline{2\pi Q_x n}$

HW: Use (A) to find

$2J_x = \beta_x x'^2 + 2\alpha_x x x' + \gamma x^2$   $\gamma \equiv \frac{1+\alpha^2}{\beta}$

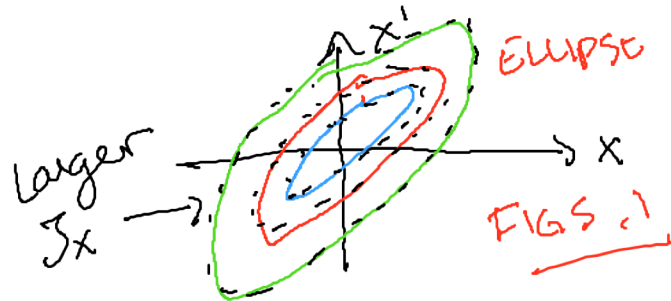
Courant-Snyder invariant  $\rightarrow$  EMITTANCE  $[J_x] = [L]$

#### 4 Ellipse and RMS displacement

$$2J_x = \beta_x \overline{x'^2} + 2\alpha_x \overline{xx'} + \gamma_x \overline{x^2}$$

↓ draw in  $(x, \phi')$

single particles



$$A_x = \text{area of ellipse} = 2\pi J_x \quad (\text{do math})$$

"single-particle emittance" (area in phase space)

True everywhere in system:  $J_x$  has NO s-dependence

RMS displacement of this single particle over all s

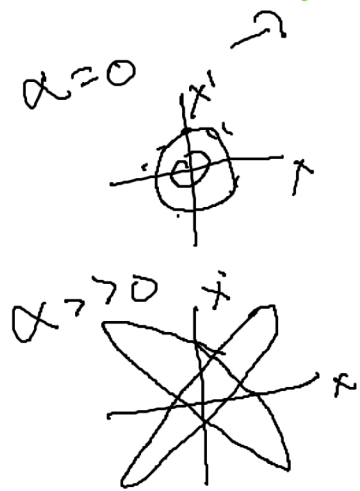
(over all  $\phi_x$ ):  $\sigma_x = \langle x^2 \rangle^{1/2} = 2\sqrt{J_x \beta_x} \langle \sin^2 \phi \rangle$

$$x = \sqrt{2J_x \beta_x} \sin \phi_x$$

$$= \sqrt{J_x \beta_x} = \sigma_x$$

"simple" 1-particle

⇒ onward to distributions



5 Many particles: distributions

Base on 1-particle statistics

Mean-square size of "bunch":  $\sigma_x^2 = \frac{1}{N} \int_0^\infty \underline{\sigma_{x_i}^2} p(J_x) dJ_x$

⇒ remember different particles have different  $J_x$

$N \propto 10^6 - 10^{11}$  or more ----

$p(J_x)$  is distribution function

$$\sigma_x^2 = \frac{\beta_x}{N} \int_0^\infty J_x p(J_x) dJ_x$$

$$\int_0^\infty p(J_x) dJ_x = \frac{1}{N} \quad (N \text{ for bunch})$$

$$\sigma_x^2 = \beta_x \langle J_x \rangle \quad \text{RMS beam size}$$

DEFINE UNNORMALIZED EMITTANCE (RMS)

$$E_{x,1} = \langle J_x \rangle$$

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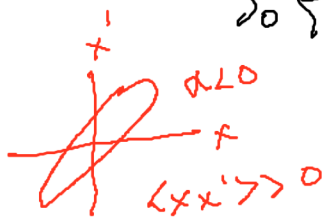
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$$\sigma_x^2 = \beta_x \langle J_x \rangle \quad \text{RMS beam size}$$



DEFINE UNNORMALIZED EMITTANCE (RMS)

$$E_{x,1} = \langle J_x \rangle$$

Generalize:

RMS quantities

$$\begin{pmatrix} \langle x^2 \rangle \\ \langle xx' \rangle \\ \langle x'^2 \rangle \end{pmatrix} = \underline{E_x} \begin{pmatrix} \beta_x \\ -\alpha_x \\ \gamma_x \end{pmatrix}$$

C-S or Twiss parameters for bunch or beam

Redid linear transverse dynamics: What about longitudinal? Dispersion?

6: Add dispersion to beam

For each particle  $x_{n, \text{TOT}} = \underbrace{\sqrt{2J_x \beta_x}}_{\text{betatron}} \sin \phi_{x,n} + \underbrace{\delta \eta_x(s)}_{\text{dispersion}}$   $\delta = \frac{\Delta p}{p_0}$

$\delta_n = \alpha_s \sin(\phi_{s,n})$

Assume  $\phi_{x,n}$  &  $\phi_{s,n}$  are uncorrelated

$$\sigma_{x, \text{TOT}}^2 = \beta_x \epsilon_x + \gamma_x \left( \frac{\sigma_p}{p} \right)^2$$

*Basic correction*  $\rightarrow 2$

ADDED "in quadrature"  
all rms quantities

(Synchrobetatron coupling is largely pathological and not common)

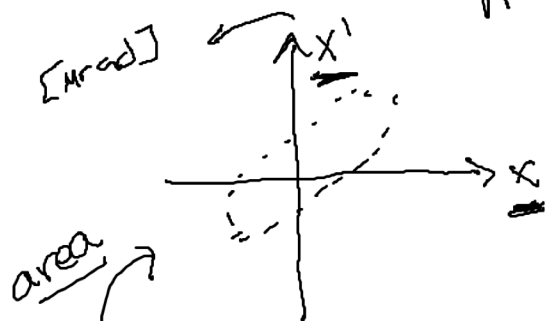
5.3: Tune spread and filamentation — website video

So far: unnormalized emittance  
 $\Rightarrow$  normalized emittance

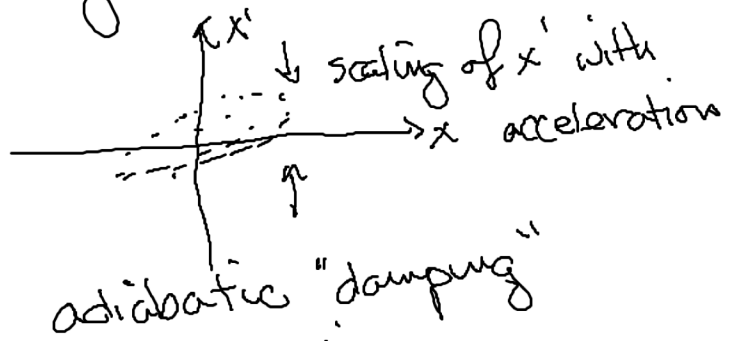


# 7: Normalized emittance and adiabatic damping

What happens to phase space during acceleration?



$$x' \equiv \frac{dx}{ds} = \frac{px}{p_0}$$



accelerate  $p_0 \rightarrow p_0 + \Delta p$

$x' \rightarrow x' / (1 + \frac{\Delta p}{p_0})$  & shrinks during acceleration

$$\langle J_{\text{new}} \rangle = \frac{J_0}{1 + \Delta p/p_0} = \frac{J_0}{(1 + \Delta(\beta\gamma)/\beta\gamma)}$$

$$p = \beta\gamma mc \quad \leftarrow \text{relativistic}$$

UNNORMALIZED EMITTANCE SHRINKS! NOT CONSERVED! (:(

Define norm emit:  $\underline{\epsilon}_{x,n} = (\beta\gamma) \epsilon_x \leftarrow \text{rms emit}$

↳ conserved even through acceleration

$$\sigma_x^2 = \frac{\beta\epsilon_{x,n}}{(\beta\gamma)}$$

$$[\epsilon_x] = [\epsilon_{x,n}] = [L]$$

Commonly  $\epsilon_x \sim \mu\text{m}$   
 $\epsilon_{x,n} \sim \text{MM-Mrad}$   
 Pion

For protons/ions: synchrotrons  $E_x \sim E_y$   $\sigma_x \sim \sigma_y$   
ROUND BEAMS

(Symmetry to transverse dynamics)

synch radiation!

$\Rightarrow$  adds  $\times$  noise (bending plane)

$\Rightarrow E_x \gg E_y$   $\sigma_x \gg \sigma_y$

FLAT BEAMS


For electrons

$$\mathcal{L} \propto \frac{I}{\sigma_x \sigma_y} \quad [cm^{-2} s^{-1}]$$

EIC:  $\Rightarrow$  Problem!

$\hookrightarrow$  electron-ion  
collider 

TLA (Three letter acronym)

p ion beams want to be round   
e beams want to be flat

Single-pass systems

LINACS

TRANSFER LINES

⇒ No periodicity so cannot use  $M = \begin{pmatrix} \cos \mu + \alpha \sin \mu & \beta \sin \mu \\ \dots & \dots \end{pmatrix}$



PARADIGM SHIFT: use  $\beta, \alpha, \gamma$  of beam distributions

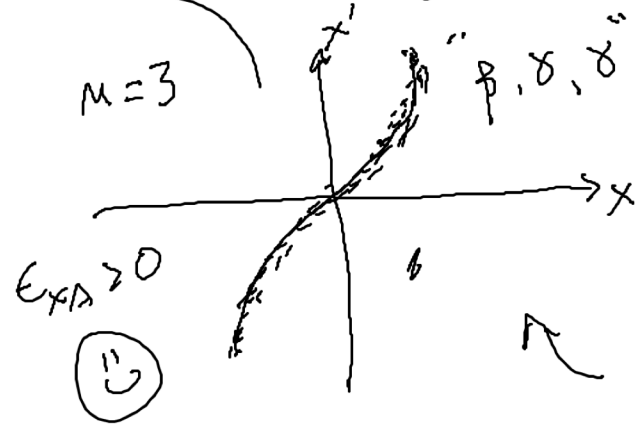
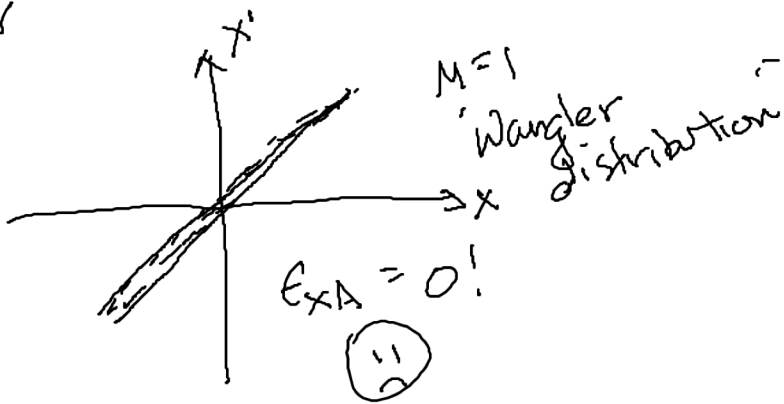
↳ This form is only for periodic focusing

electrons  
↓

$$\begin{pmatrix} \beta \\ -\alpha \\ \gamma \end{pmatrix}_{x'} = \frac{1}{E_{xA}} \begin{pmatrix} \langle x^2 \rangle \\ \langle xx' \rangle \\ \langle x'^2 \rangle \end{pmatrix}$$

$E_{xA} \sim$  AREA EMITTANCE (scaling term)

USES  $e^-$



$e^-$  guns/sources  
low  $E$  transport  
non-equilibrium systems  
(CEZLS, FELS...)

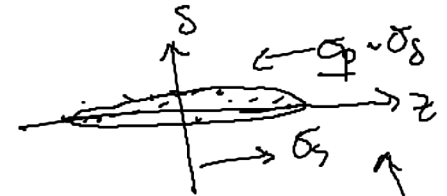
Are there longitudinal equivalents for this description?

longitudinally RMS bunch length  $\sigma_s \equiv \sqrt{\frac{\beta_s E_s}{(\beta \gamma)^2}}$

normalization  $\rightarrow$

$$\sigma_p = \sqrt{\frac{E_s}{\beta_s (\beta \gamma)^2}}$$

$\beta_s$ : longitudinal  $\beta$  function parameterization



$E_s \equiv$  longitudinal emittance  
 units could be [L]  $\Leftrightarrow z, s$   
 often see  $\frac{eV \cdot s}{\text{time}}$   
energy spread

$\beta_s [M]$  =  $\frac{\sigma_s}{\sigma_p}$  = easily calculable

circumf  $\rightarrow \frac{C}{2\pi}$

$\frac{|H_s|}{Q_s} \leftarrow$  slip factor =  $-\frac{1}{\gamma^2} + \frac{1}{\gamma_j^2}$   
 $\leftarrow$  synchrotron tune

when  $\delta = \delta_T$   
 $\mu_s = 0$

$\sigma_p \sim \frac{1}{\sqrt{\beta_s}} \rightarrow \infty!$   
 $\rightarrow \beta_s \rightarrow 0$   
 BOOM 😊