Beam Cooling

M. Steck, GSI, Darmstadt

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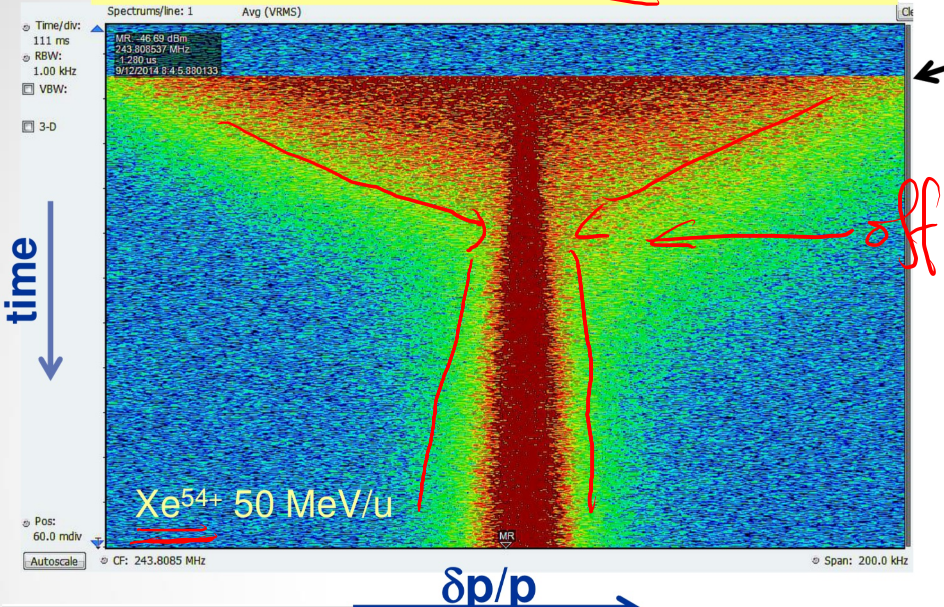
Beam Cooling ⇒ Reducing ϵ

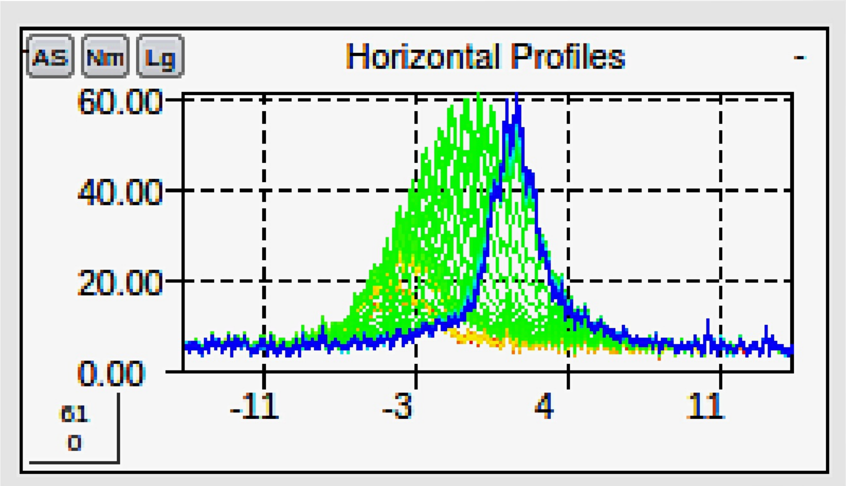
↳ Stochastic cooling Nobel!

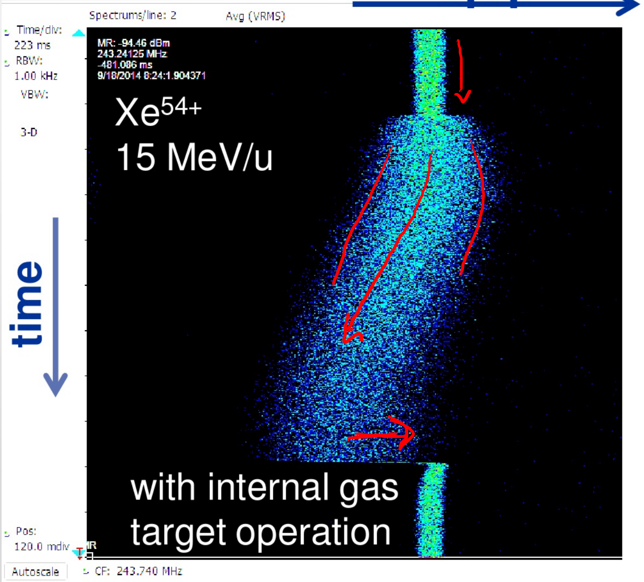
longitudinal (momentum) cooling

injection into the storage ring



transverse cooling





cooling **off**

⇐

heating (spread) and energy loss (shift)

cooling **on**

cooling:
good energy definition
small beam size
⇒ highest precision

Beam Cooling

Introduction

1. Electron Cooling
2. Ionization Cooling
3. Laser Cooling
4. Stochastic Cooling

Beam Cooling

→ Thermodynamics

Beam cooling is synonymous for a reduction of beam temperature.

Temperature is equivalent to terms as phase space volume, emittance and momentum spread. → emittance

Beam Cooling processes are not following Liouville‘s Theorem:

`in a system where the particle motion is controlled by external conservative forces the phase space density is conserved´

(This neglects interactions between beam particles.)

Beam cooling techniques are non-Liouvillean processes which violate the assumption of a conservative force.

e.g. interaction of the beam particles with other particles
(electrons, photons, matter)

Cooling Force

Generic (simplest case of a) cooling force: Friction

$$F_{x,y,s} = -\underline{\alpha_{x,y,s}}\,\underline{v_{x,y,s}}$$

$v_{x,y,s}$ velocity in the rest frame of the beam

non conservative, cannot be described by a Hamiltonian

Dissapative

For a 2D subspace distribution function $f(z, z', t)$

$$F_z = -\alpha_z v_z \qquad z = x, y, s \qquad v_z = v_0 \cdot z'$$

$$\frac{df(z, z', t)}{dt} = -\underline{\lambda_z f(z, z', t)}$$

$\underline{\lambda_z}$ cooling (damping) rate $\propto$ 1/cooling time

in a circular accelerator:

Transverse (emittance) cooling

$$\epsilon_{x,y}(t_0 + t) = \epsilon_{x,y}(t_0)\; \underline{e^{-\lambda_{x,y} t}}$$

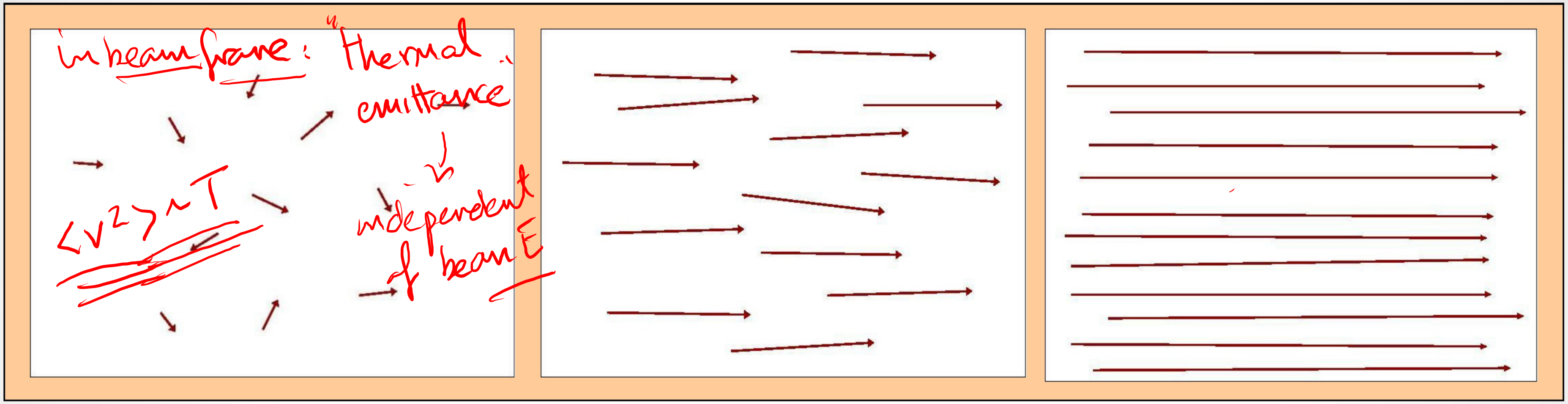
Longitudinal (momentum spread) cooling

$$\frac{\delta p_{\|}}{p_0}(t_0 + t) = \frac{\delta p_{\|}}{p_0}(t_0)\; \underline{e^{-\lambda_{\|} t}}$$

Beam Temperature

Where does the beam temperature originate from?

The beam particles are generated in a 'hot' source



in beam frame: "thermal emittance"

→ independent of beam E

$\langle v^2 \rangle \sim T$

at rest (source) **at low energy** **at high energy**

In a standard accelerator the beam temperature is not reduced (thermal motion is superimposed the average motion after acceleration)

but: many processes can heat up the beam → internal & conservative

e.g. heating by mismatch, space charge, intrabeam scattering, internal targets, residual gas, external noise

Beam Temperature Definition

Longitudinal beam temperature

$$\boxed{\underline{\frac{1}{2}k_BT_\parallel} = \underline{\left\langle\frac{1}{2}mv_\parallel^2\right\rangle} = \frac{1}{2}mc^2\beta^2\left(\frac{\delta p_\parallel}{p}\right)^2}$$

Transverse beam temperature

$$\boxed{\underline{\frac{1}{2}k_BT_\perp} = \underline{\left\langle\frac{1}{2}mv_\perp^2\right\rangle} = \frac{1}{2}mc^2\beta^2\gamma^2\theta_\perp^2}$$

$$\theta_\perp = \frac{v_\perp}{\beta c}, \qquad \theta_\perp(s) = \sqrt{\frac{\epsilon}{\beta_\perp(s)}}$$

Distribution function

$$f(v_\perp, v_\parallel) \propto \exp\left(-\frac{mv_\perp^2}{2k_BT_\perp} - \frac{mv_\parallel^2}{2k_BT_\parallel}\right)$$

Particle beams can be anisotropic: $\underline{k_BT_\parallel} \neq \underline{k_BT_\perp}$ ⇐

e.g. due to laser cooling or the distribution of the electron beam

↪ Don‘t confuse: beam energy ↔ beam temperature

(e.g. a beam of energy 100 GeV can have a temperature of 1 eV)

Benefits of Beam Cooling

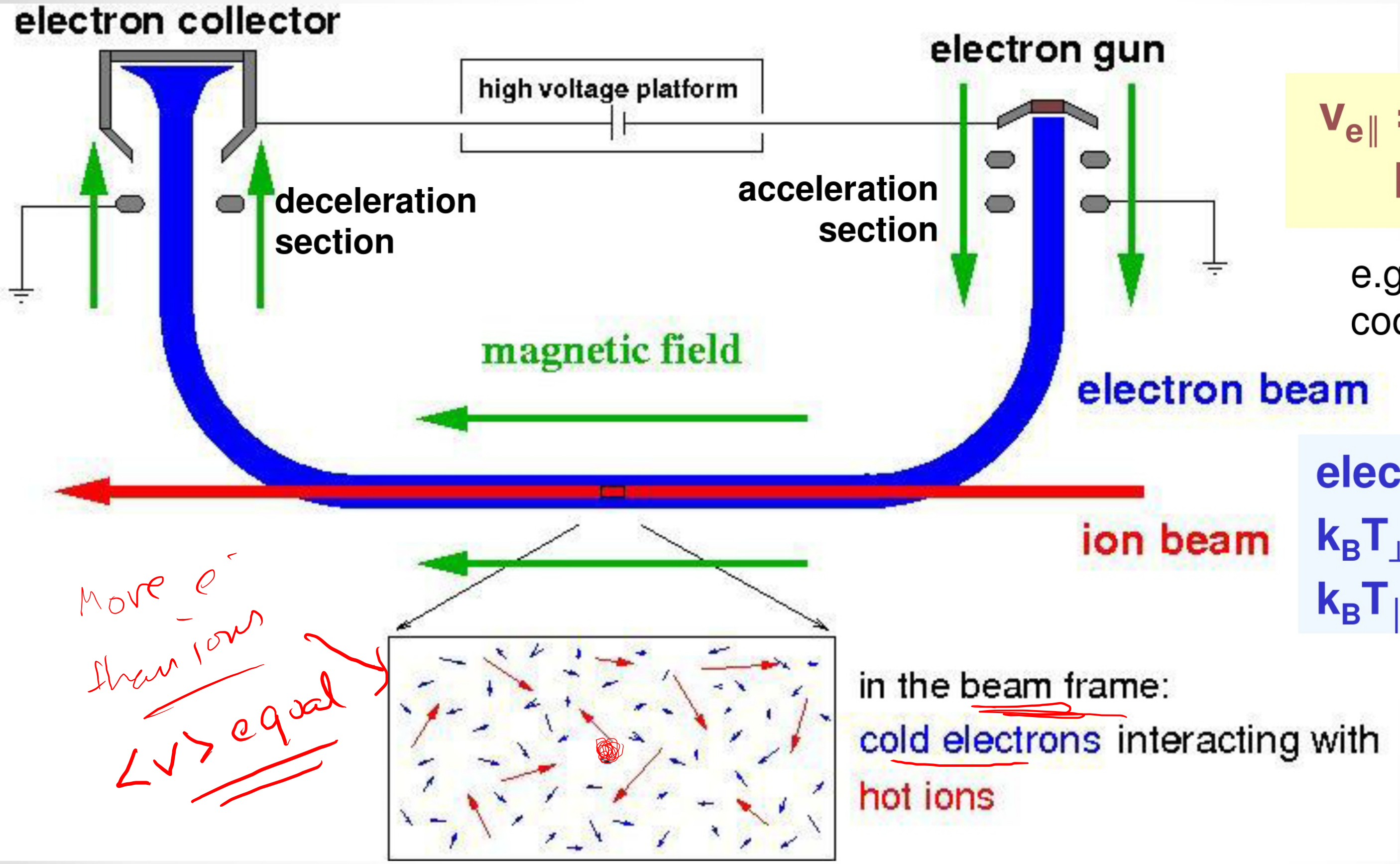
- Improved beam quality
  - Precision experiments
  - Luminosity increase

→ EIC will rely on electron cooling

- Compensation of heating
  - Experiments with internal target
  - Colliding beams
- Intensity increase by accumulation
  - Weak beams from the source can be enhanced
  - Secondary beams (antiprotons, rare isotopes)

1. Electron Cooling ⇒

cold gas → e^- beam
mix hot gas → hadron beam



in the beam frame:
cold electrons interacting with
hot ions

more e^- than ions
$\langle v \rangle$ equal

$$v_{e\parallel} = \beta_e c = \beta_i c = v_{i\parallel}$$
$$E_e = m_e / M_i \cdot E_i$$

e.g.: 220 keV electrons
cool 400 MeV protons

electron temperature
$k_B T_\perp \approx 0.1$ eV
$k_B T_\parallel \approx 0.1 - 1$ meV

thermal transfer

superposition of a cold intense electron beam with the same velocity

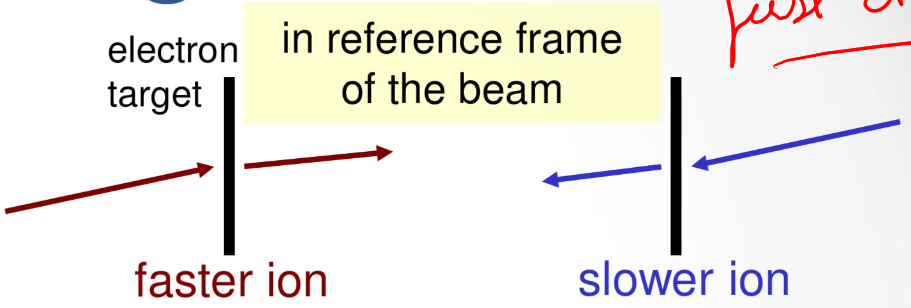
momentum transfer by Coulomb collisions

cooling force results from energy loss in the co-moving gas of free electrons

Simple Derivation of the Electron Cooling Force

⇒ Friction to first order

Analogy: energy loss in matter (electrons in the shell)


electron target

in reference frame of the beam

faster ion

slower ion

Rutherford scattering: $2\ tan(\frac{\theta}{2}) = \frac{2Z_1Z_2e^2}{4\pi\epsilon_0\Delta p v b} \quad Z_1 = Q\ (ion),\ \ Z_2 = -1\ (electron)$

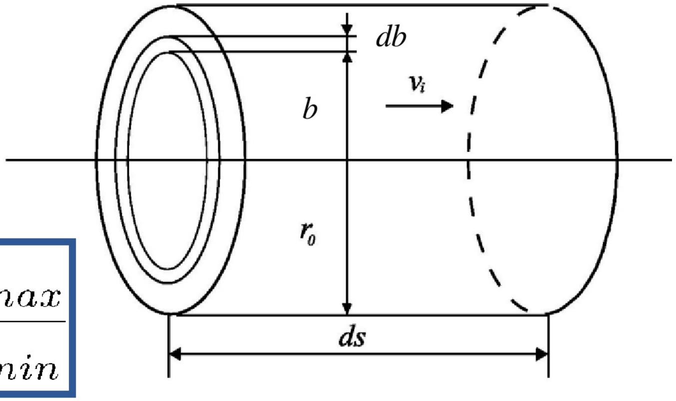
Energy transfer: $\Delta E(b) = \frac{(\Delta p)^2}{2m_e} \simeq \frac{2Q^2e^4}{(4\pi\epsilon_0)^2 m_e v^2}\frac{1}{b^2} \quad (for\ \ b \gg b_{min})$

Minimum impact parameter: $b_{min} = \frac{Qe^2}{(4\pi\epsilon_0)^2 m_e v^2}$

from: $\Delta E(b_{min}) = \Delta E_{max} \simeq 2m_e v^2$

Energy loss:

$$\boxed{-\frac{dE}{dx} = 2\pi\int_{b_{min}}^{b_{max}} b n_e \Delta E\, db = \frac{4\pi Q^2e^4}{(4\pi\epsilon_0)^2 m_e v^2} n_e \ln\frac{b_{max}}{b_{min}}}$$



Coulomb logarithm $L_C = \ln(b_{max}/b_{min}) \approx 10$ (typical value)

Characteristics of the Electron Cooling Force

$$\boxed{\vec{F}(\vec{v_i}) = -\frac{4\pi Q^2 e^4 n_e}{(4\pi\epsilon_0)^2 m_e}\int L_C(\vec{v}_{rel}) f(\vec{v_e}) \frac{\vec{v}_{rel}}{v_{rel}^3} d^3\vec{v_e}}$$

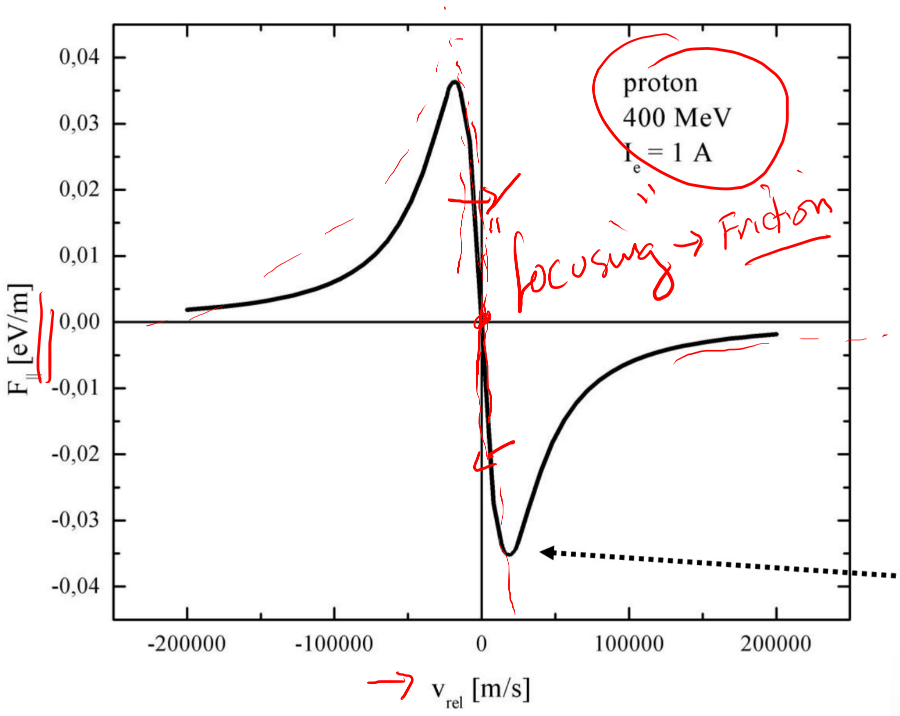
$$\vec{v}_{rel} = \vec{v_i} - \vec{v_e}$$

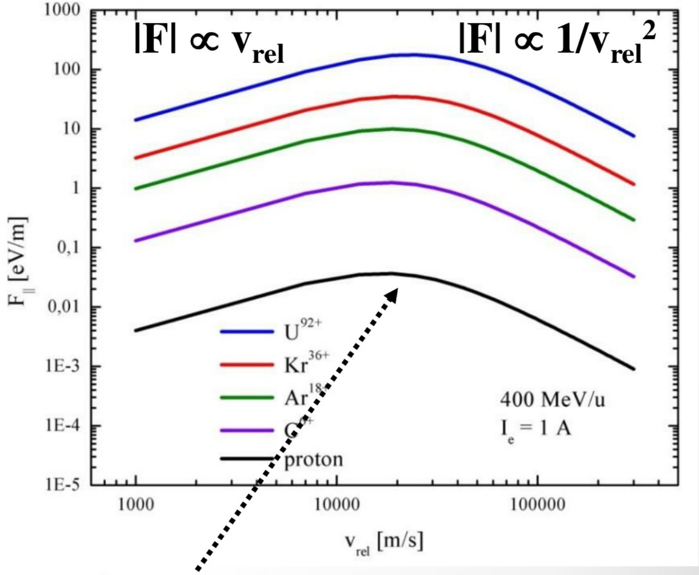
cooling force F → beam-beam space charge

for small relative velocity: $\propto v_{rel}$

for large relative velocity: $\propto v_{rel}^{-2}$

increases with charge: $\propto Q^2$ → ions





maximum of cooling force
at effective electron temperature

Electron Cooling Time

first estimate:
(Budker 1967)

$$\tau = \frac{3}{8\sqrt{2\pi}n_e Q^2 r_e r_i c L_C}\left(\frac{k_B T_e}{m_e c^2} + \frac{k_B T_i}{m_i c^2}\right)^{3/2}$$

for large relative velocities

relativistic β, γ

cooling time $\tau_z \propto \frac{A}{Q^2}\frac{1}{n_e \eta}\beta^3\gamma^5\theta_z^3$ $\begin{cases} \theta_{x,y} = \frac{v_{x,y}}{\gamma\beta c} \\ \theta_{\parallel} = \frac{v_{\parallel}}{\gamma\beta c} \end{cases}$

βγθ ~ conserved

extra $\gamma^2 \propto \tau_z$

↳ Loses effectiveness as γ^{-2}

Time dilation ←

cooling rate (τ⁻¹):

- slow for hot beams $\propto \theta^{-3}$
- decreases with energy $\propto \gamma^{-2}$ (β·γ·θ is conserved)
- linear dependence on electron beam intensity n_e and cooler length $\eta = L_{ec}/C$
- favorable for highly charged ions Q^2/A
- independent of hadron beam intensity ←

for small relative velocities

cooling rate is constant and maximum at small relative velocity

$F \propto v_{rel} \Rightarrow \tau = \Delta t = p_{rel}/F = \text{constant}$

Models of the Electron Cooling Force

- binary collision model

description of the cooling process by successive collisions of two particles and integration over all interactions
analytic expressions become very involved, various regimes (multitude of Coulomb logarithms)

- dielectric model

interaction of the ion with a continuous electron plasma (scattering off of plasma waves)
fails for small relative velocities and high ion charge

- an empiric formula (Parkhomchuk) derived from experiments:

$$\boxed{\vec{F} = -4\frac{n_e}{m_e}\frac{(Qe^2)^2}{(4\pi\epsilon_0)^2}\ln\left(\frac{b_{max}+b_{min}+r_c}{b_{min}+r_c}\right)\frac{\vec{v}_{ion}}{(v_{ion}^2+v_{eff}^2)^{3/2}}}$$

$$b_{min} = \frac{Qe^2/4\pi\epsilon_0}{m_e v_{ion}^2}; \quad b_{max} = \frac{v_{ion}}{min(\omega_{pe}, 1/T_{cool})}, \quad v_{eff}^2 = v_{e,\parallel}^2 + v_{e,\perp}^2$$

Electron Beam Properties

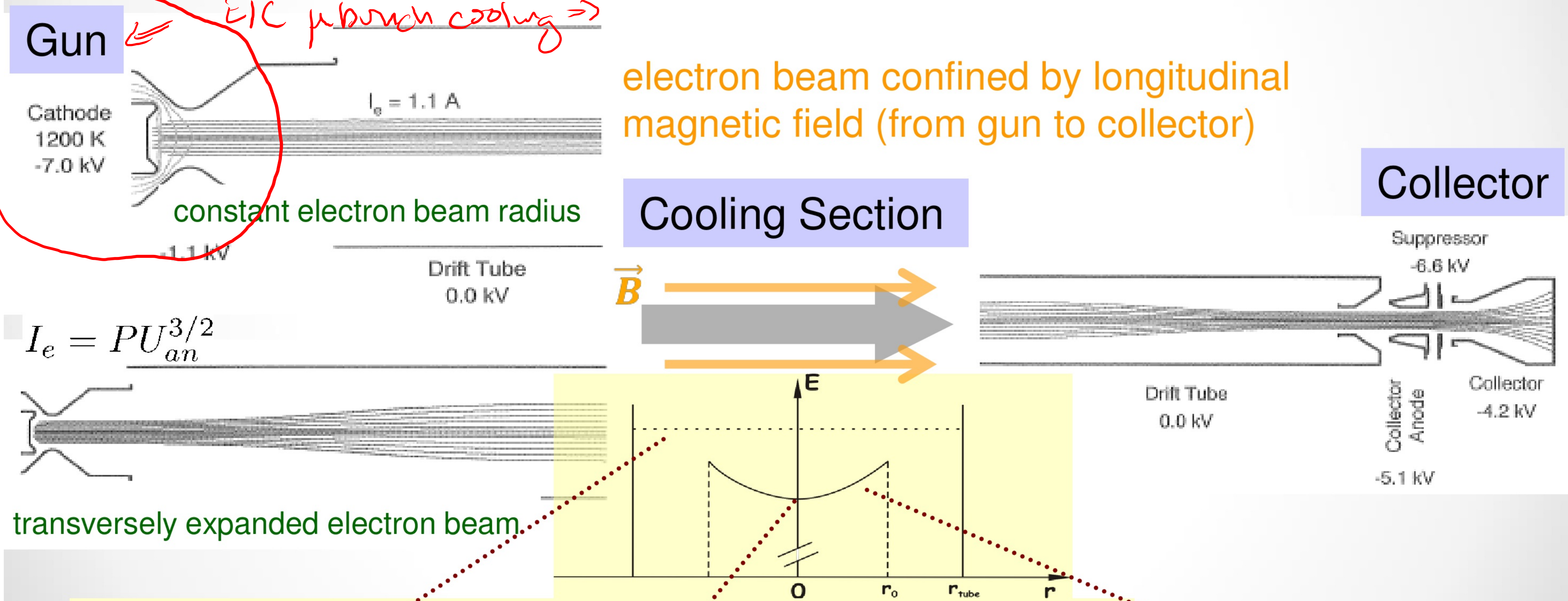
electron beam temperature

transverse $k_B T_\perp = k_B T_{cat}$, with transverse expansion ($\propto B_c/B_{gun}$)

longitudinal $k_B T_\parallel = (k_B T_{cat})^2/4E_0 \ll k_B T_\perp$ lower limit : $k_B T_\parallel \geq 2e \frac{n_e^{1/3}}{4\pi\epsilon_0}$

typical values: $k_B T_\perp \approx 0.1$ eV (1100 K), $k_B T_\parallel \approx 0.1 - 1$ meV

EIC μbunch cooling =>



electron beam confined by longitudinal magnetic field (from gun to collector)

constant electron beam radius

$$I_e = PU_{an}^{3/2}$$

transversely expanded electron beam

radial variation of electron energy due to space charge:

$$E(r) = eU_{cat} - \underline{n_e}\pi r_0^2 r_e m_e c^2[1 + 2\ln(r_{tube}/r_0)] + \underline{n_e}\pi r_e m_e c^2 r^2$$

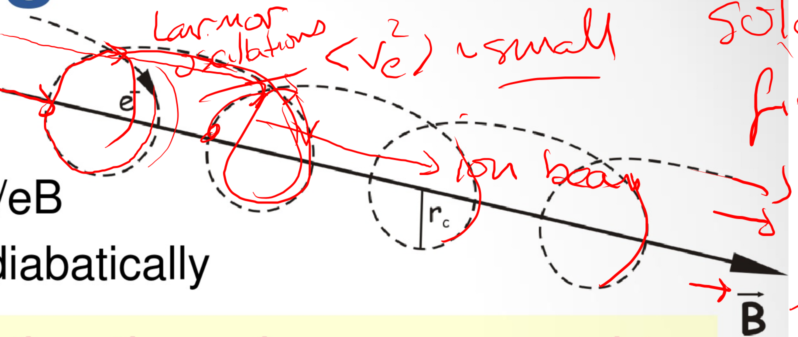
Electron Motion in Longitudinal Magnetic Field

single particle cyclotron motion

cyclotron frequency $\omega_c = eB/\gamma m_e$

cyclotron radius $r_c = v_\perp/\omega_c = (k_B T_\perp m_e)^{1/2}\,\gamma/eB$

electrons follow the magnetic field line adiabatically



solenoidal field &

ion divergence small

⇓

$\alpha_{x,y}$ small

$\beta_{x,y}$ constant

β_{xy} large & constant

important consequence: for interaction times long compared to the cyclotron period the ion does not sense the transverse electron temperature ⇒ **magnetized cooling ($T_{eff} \approx T_\parallel << T_\perp$)**

electron beam space charge:

transverse electric field + B-field ⇒ azimuthal drift

$$v_{azi} = r\omega_{azi} = r\frac{2\pi r_e n_e c^2}{\gamma\omega_c}$$

⇒ electron and ion beam should be well centered

Favorable for optimum cooling (small transverse relative velocity):

- high parallelism of magnetic field lines $\Delta B_\perp/B_0$
- large beta function (small divergence) in cooling section

Imperfections and Limiting Effects in Electron Cooling

technical issues:

ripple of accelerating voltage

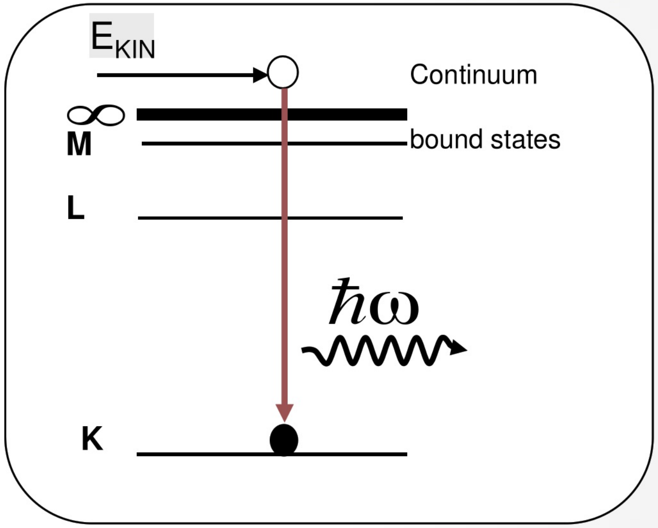
magnetic field imperfections

beam misalignment

space charge of electron beam and compensation

physical limitation:

Radiative Electron Capture (REC)



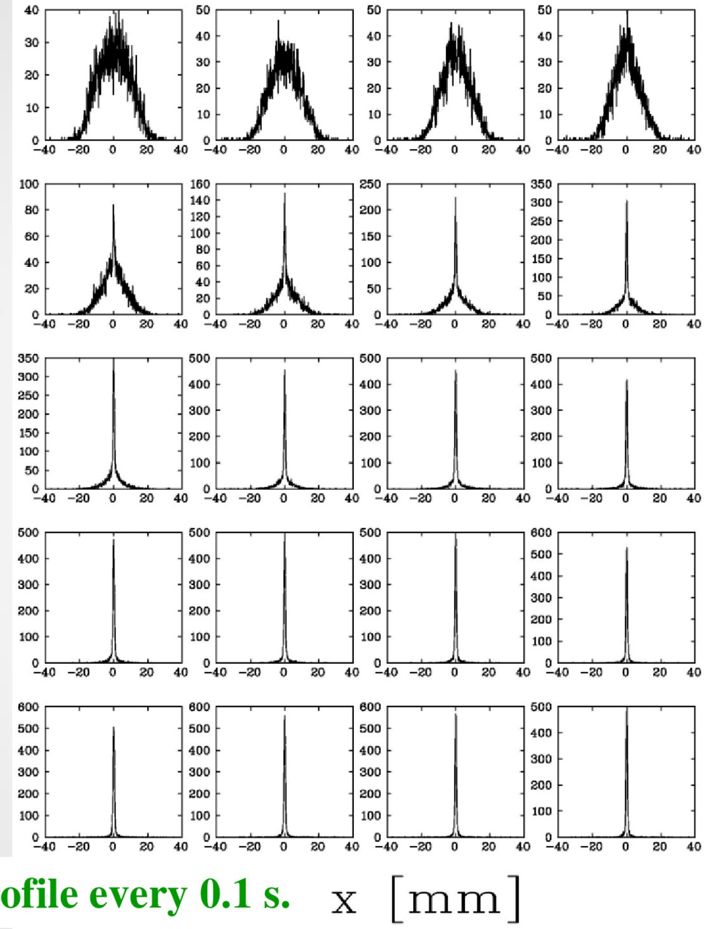
losses by recombination (REC)

loss rate $\tau^{-1} = \gamma^{-2} \alpha_{REC} n_e \eta$

$$\alpha_{REC} = \frac{1.92 \times 10^{-13} Q^2}{\sqrt{k_B T}} \left(\ln \frac{5.66 Q}{\sqrt{k_B T}} + 0.196 (\frac{k_B T}{Q^2})^{1/3} \right) [cm^3 s^{-1}]$$

Examples of Electron Cooling

fast transverse cooling at TSR, Heidelberg



profile every 0.1 s.

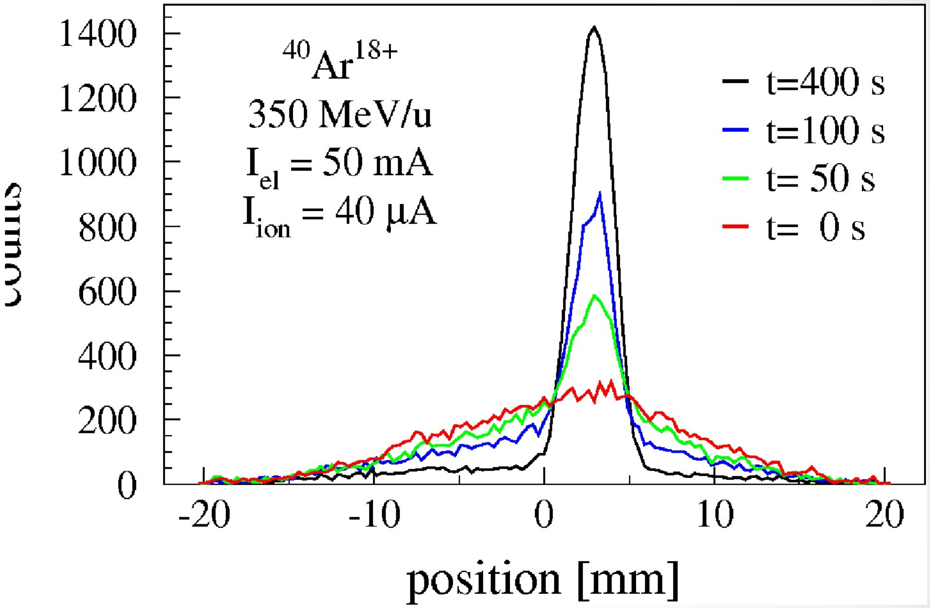
cooling of 6.1 MeV/u C^{6+} ions

0.24 A, 3.4 keV electron beam

$n_e = 1.56 \times 10^7$ cm^{-3}

measured with residual gas ionization beam profile monitor

transverse cooling at ESR, Darmstadt



cooling of 350 MeV/u Ar^{18+} ions

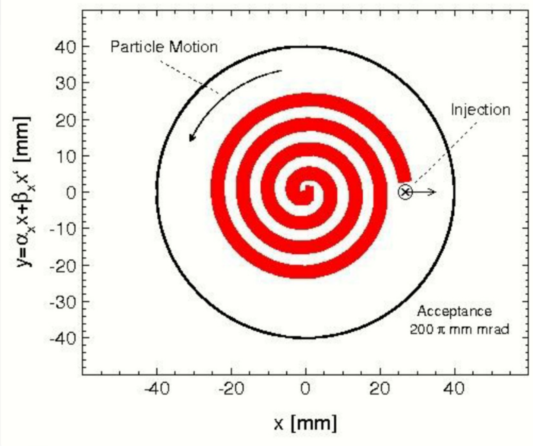
0.05 A, 192 keV electron beam

$n_e = 0.8 \times 10^6$ cm^{-3}

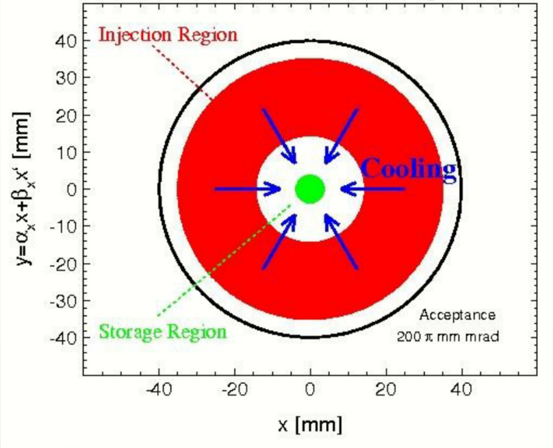
note! time scale, the cooling time varies strongly with beam parameters

Accumulation of Heavy Ions by Electron Cooling

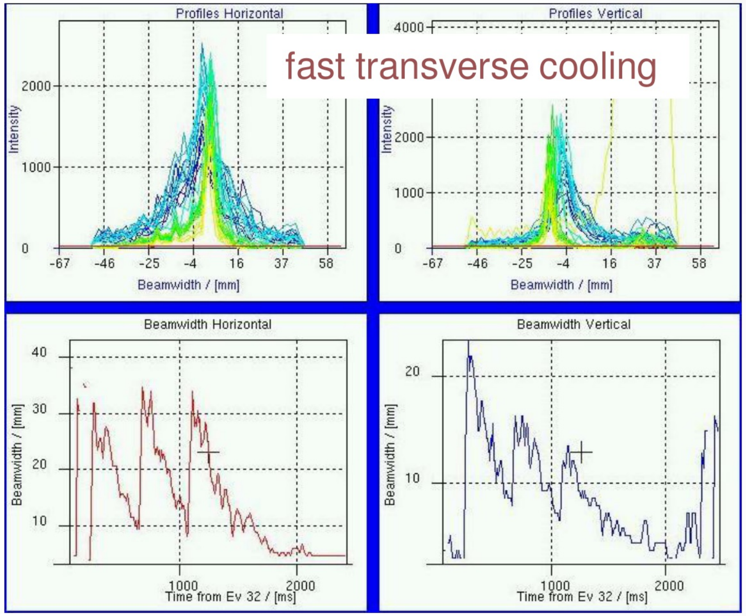
standard multiturn injection



fast accumulation by repeated multiturn injection with electron cooling



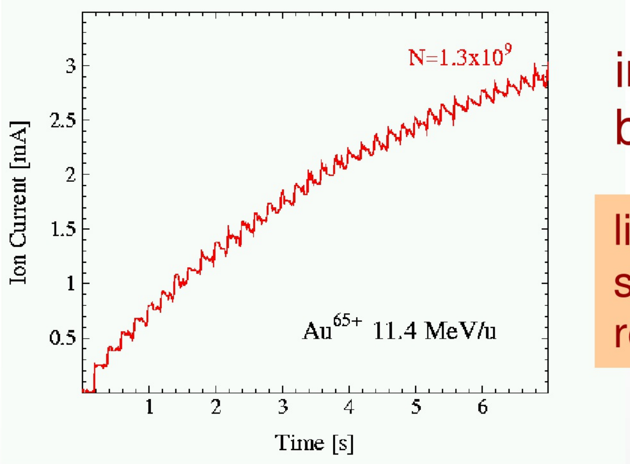
horizontal **vertical**



fast transverse cooling

profile

beam size



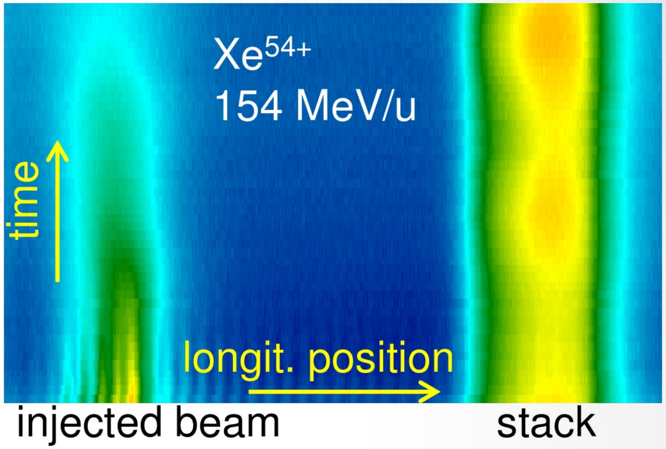
intensity increase in 5 s by a factor of ≈ 10

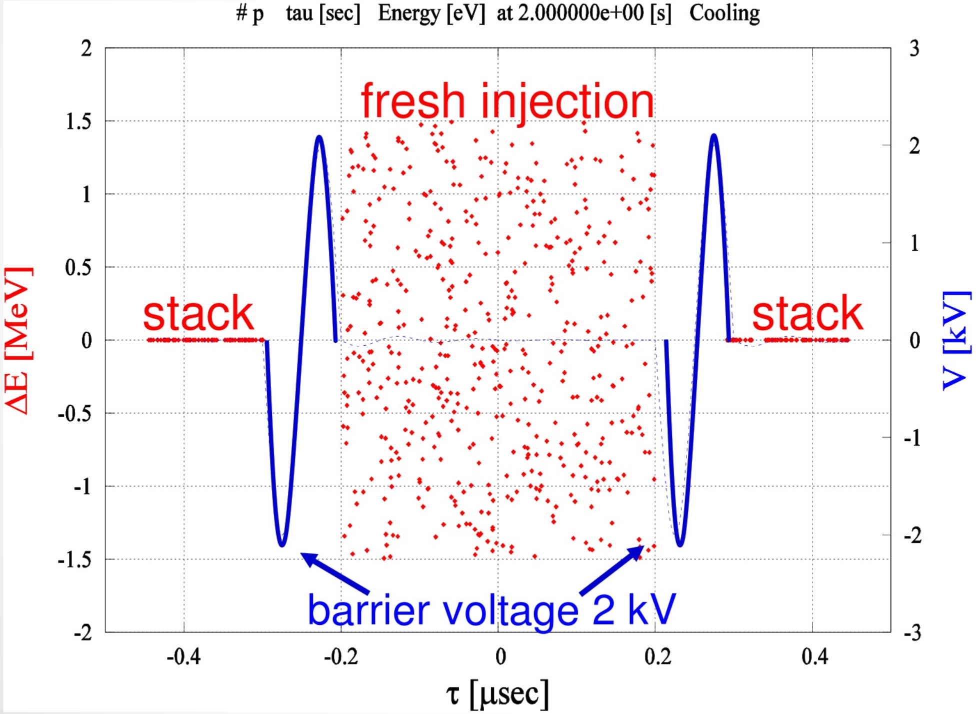
limitations:
space charge tune shift,
recombination (REC)

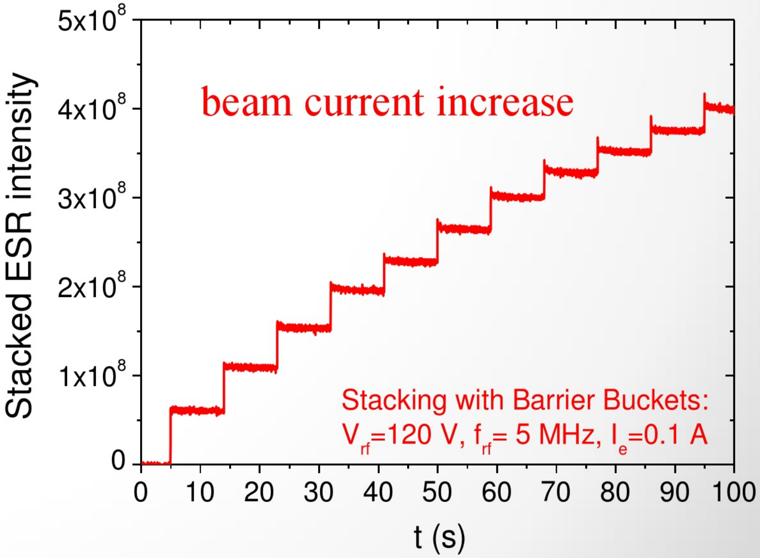
Accumulation of Secondary Particles

basic idea: confine stored beam to a fraction of the circumference, inject into gap and apply cooling to merge the two beam components
⇒ fast increase of intensity (for secondary beams)

experimental verification at ESR







Stacking with Barrier Buckets: V_{rf}=120 V, f_{rf}= 5 MHz, I_e=0.1 A

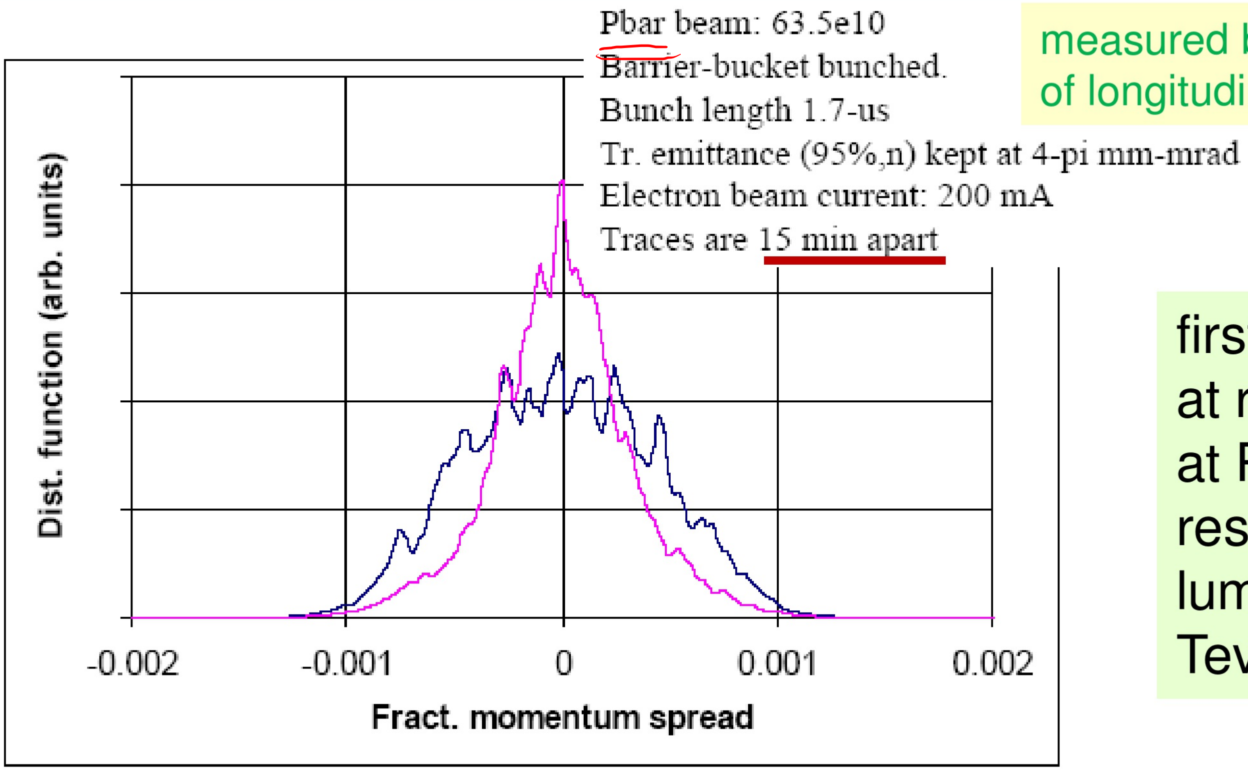
simulation of longitudinal stacking with barrier buckets and electron cooling

Examples of Electron Cooling

high energy electron cooling of 8 GeV antiprotons

longitudinal cooling with 0.2 A, 4.4 MeV electron beam

First e-cooling demonstration - 07/15/05



Pbar beam: 63.5e10
Barrier-bucket bunched.
Bunch length 1.7-us
Tr. emittance (95%,n) kept at 4-pi mm-mrad
Electron beam current: 200 mA
Traces are 15 min apart

measured by detection of longitudinal Schottky noise

first electron cooling at relativistic energy at Recycler, FNAL resulting in increased luminosity in the Tevatron collider

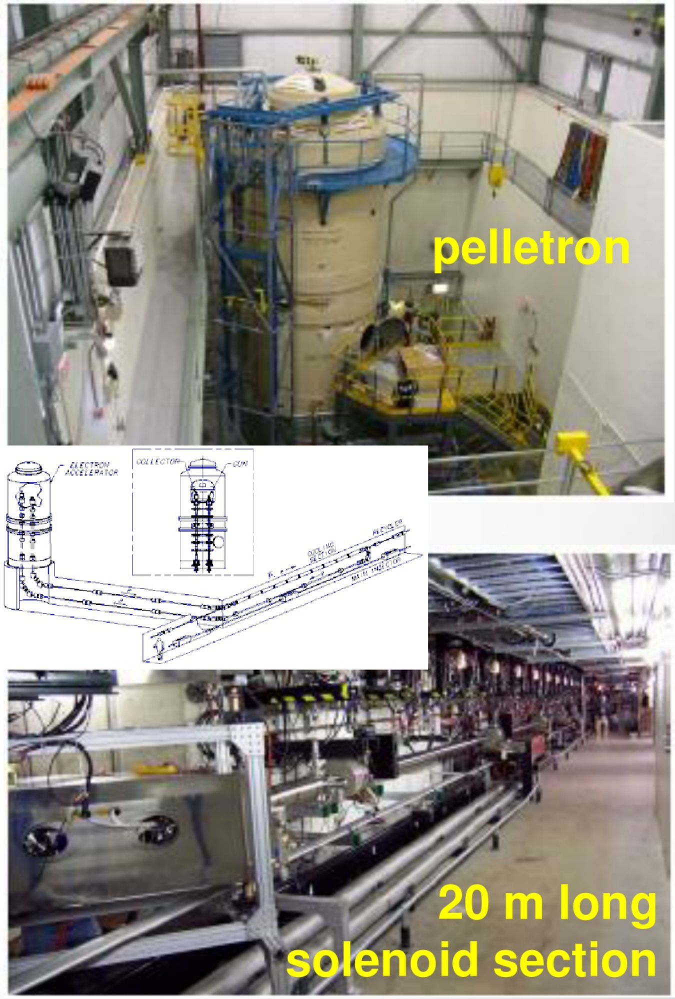
cooling time of some ten minutes has to be compared with the accumulation time of many hours

Electron Cooling Systems

Low Energy: 35 keV SIS/GSI

cooling section 3.4 m long

High Energy:
4.3 MeV Recycler/FNAL



pelletron

20 m long solenoid section

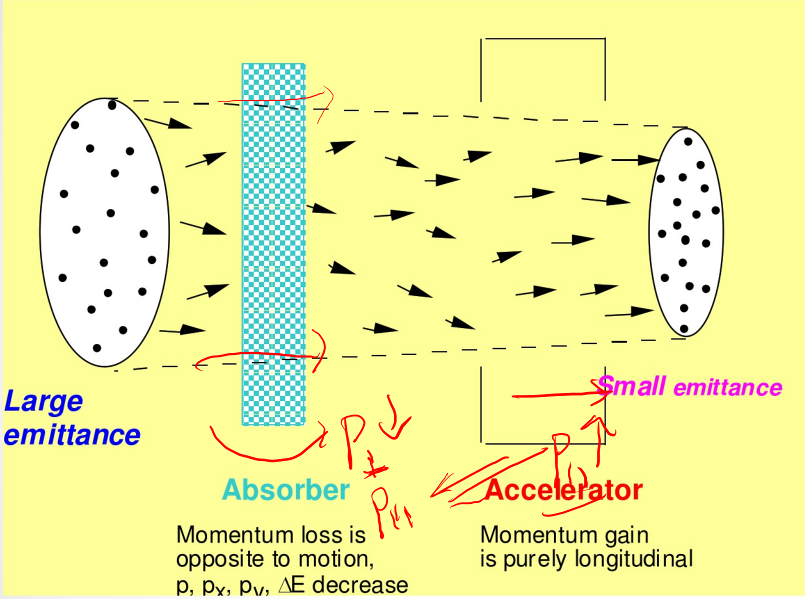
Medium Energy:
300 keV
ESR/GSI

2. Ionization Cooling

MICE expt
↳ Muon Ionization cooling experiment

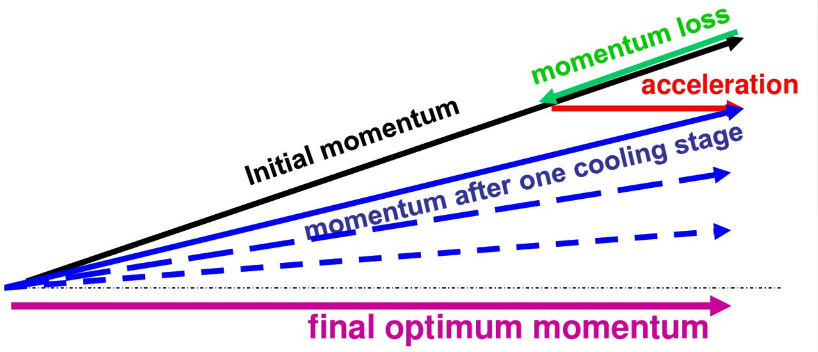
energy loss in solid matter

proposed for muon cooling



Momentum loss is opposite to motion, p, p_x, p_y, ΔE decrease

Momentum gain is purely longitudinal



not useful for heavy particles due to strong interaction with matter

transverse cooling

$$\frac{d\epsilon_N}{ds} = -\frac{1}{\beta^2 E}\frac{dE}{ds}\epsilon_N + \frac{\beta\gamma\beta_\perp}{2}\frac{\langle\theta^2_{rms}\rangle}{ds}$$

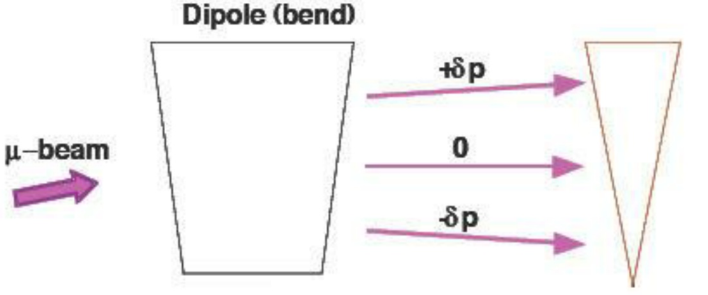
$$= -\frac{1}{\beta^2 E}\frac{dE}{ds}\epsilon_N + \frac{\beta_\perp E_s^2}{2\beta^3 m_\mu c^2 L_R E}$$

⇒ small $\beta_\perp$ at absorber in order to minimize multiple scattering

large L_R, (dE/ds) ⇒ light absorbers (H_2)

Ionization Cooling

increased longitudinal cooling
by longitudinal-transverse emittance exchange



Dipole introduces dispersion (η) $x \to x_0 + D\,\delta p/p$

Wedge Absorber reduces energy spread

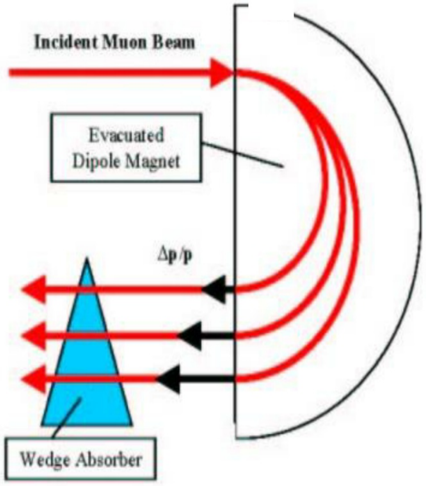


Figure 1. Use of a Wedge Absorber for Emittance Exchange

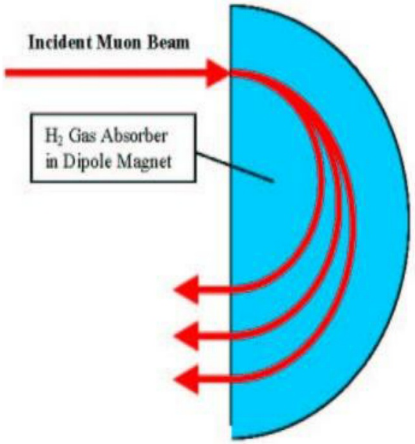


Figure 2. Use of Continuous Gaseous Absorber for Emittance Exchange

$$\frac{d\sigma_E^2}{ds} = \underbrace{-2\frac{\partial(dE/ds)}{\partial E}\sigma_E^2}_{\text{cooling term}} + \underbrace{\frac{d\langle\Delta E_{rms}^2\rangle}{ds}}_{\text{heating term}}$$

cooling, if $\frac{\partial(dE/ds)}{\partial E} > 0$

emittance exchange

increased longitudinal cooling

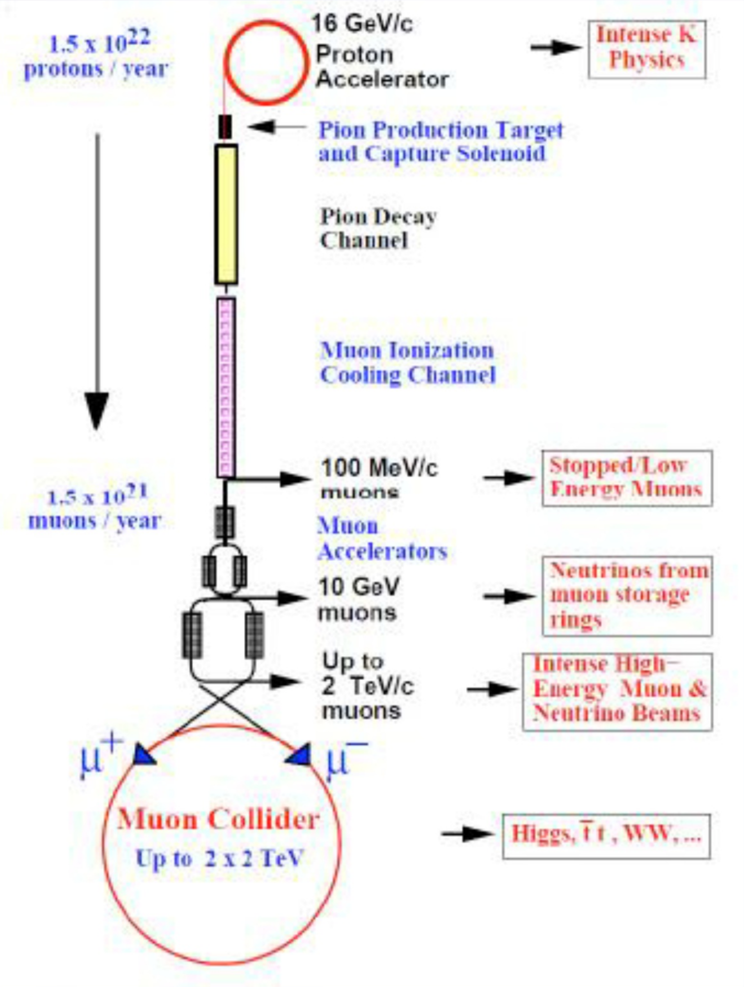
$$\frac{\partial\frac{dE}{ds}}{\partial E} \Rightarrow \frac{\partial\frac{dE}{ds}}{\partial E}\Big|_0 + \frac{dE}{ds}\frac{D\rho'}{\beta c p \rho_0}$$

reduced transverse cooling

$$\frac{d\epsilon_N}{ds} = -\frac{1}{\beta^2 E}\frac{dE}{ds}\left(1 - \frac{D\rho'}{\rho_0}\right)\epsilon_N$$

Scenarios with Ionization Cooling

Muon Collider



1.5 x 10^{22} protons / year

16 GeV/c Proton Accelerator → Intense K Physics

Pion Production Target and Capture Solenoid

Pion Decay Channel

Muon Ionization Cooling Channel

1.5 x 10^{21} muons / year

100 MeV/c muons → Stopped/Low Energy Muons

Muon Accelerators

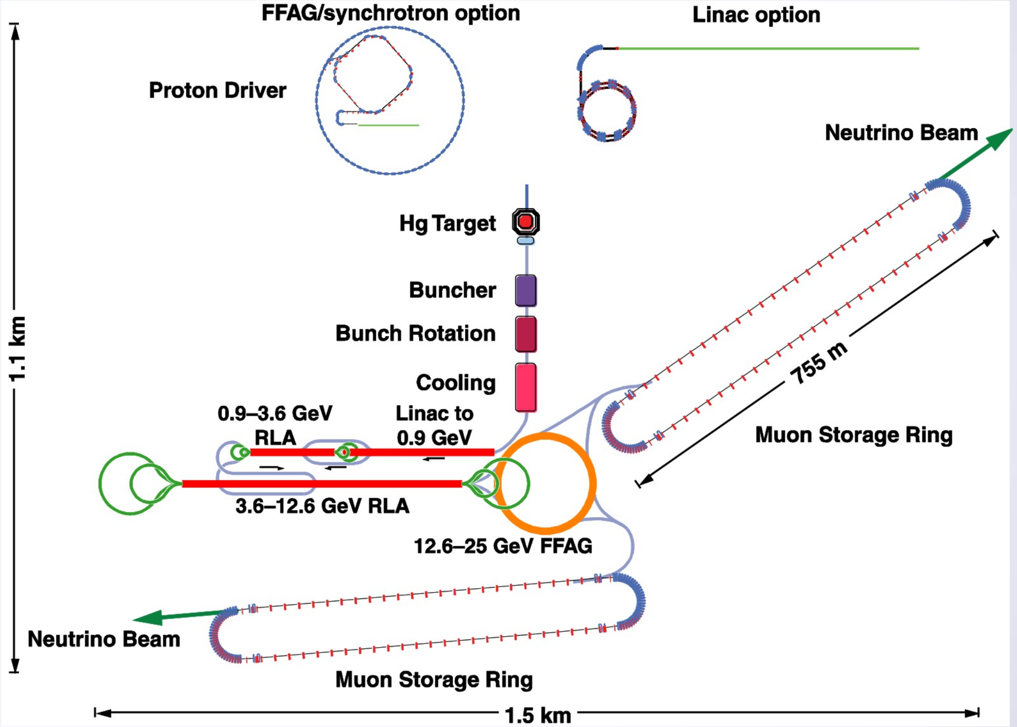
10 GeV muons → Neutrinos from muon storage rings

Up to 2 TeV/c muons → Intense High–Energy Muon & Neutrino Beams

μ^+ μ^-

Muon Collider Up to 2 x 2 TeV → Higgs, $\bar{t}t$, WW, ...

Neutrino Factory



FFAG/synchrotron option

Linac option

Proton Driver

Neutrino Beam

Hg Target

Buncher

Bunch Rotation

Cooling

0.9–3.6 GeV RLA

Linac to 0.9 GeV

3.6–12.6 GeV RLA

12.6–25 GeV FFAG

755 m

Muon Storage Ring

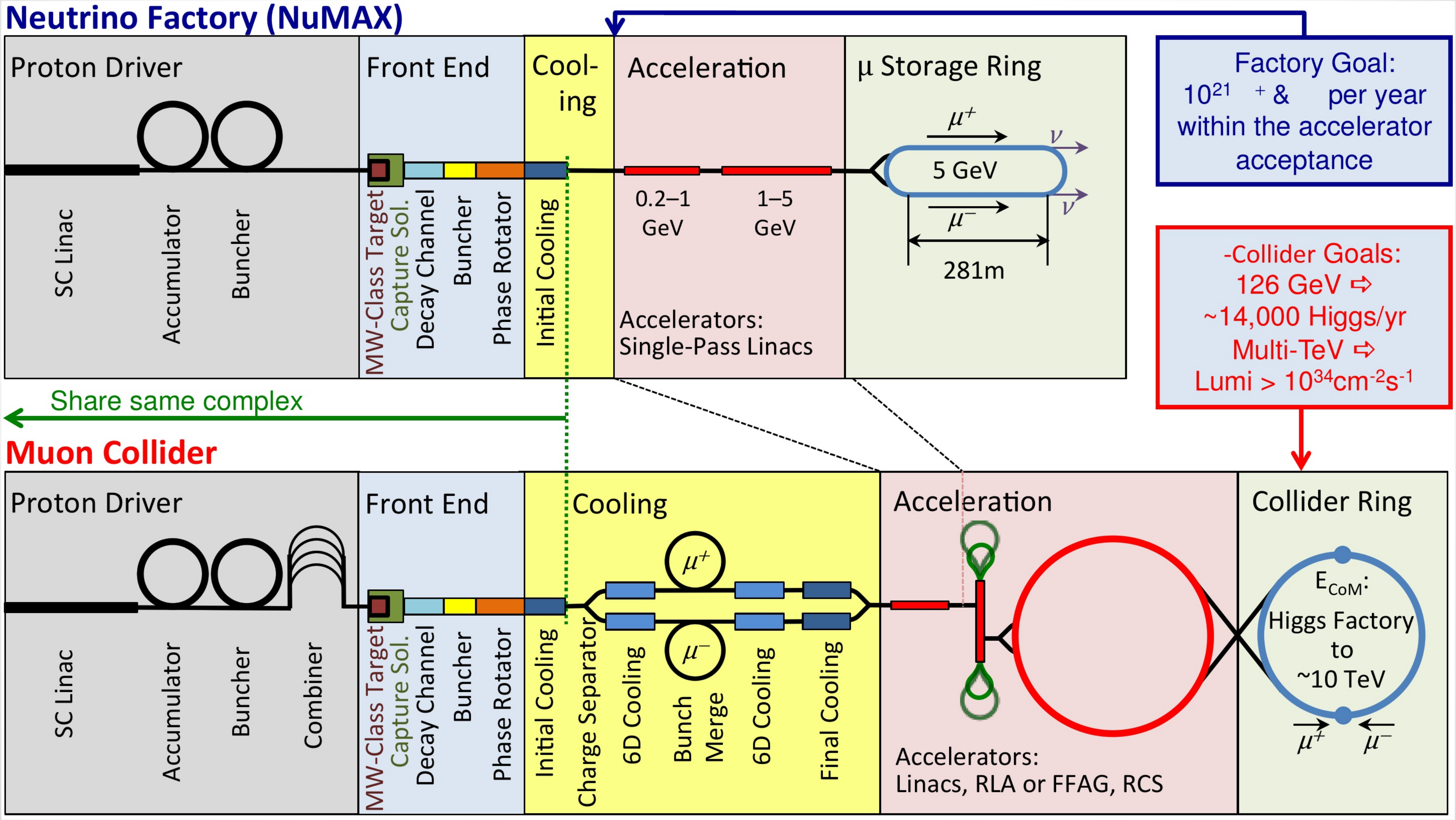
1.1 km

Neutrino Beam

Muon Storage Ring

1.5 km

Scenarios with Ionization Cooling



Neutrino Factory (NuMAX)

Proton Driver

SC Linac

Accumulator

Buncher

Front End

MW-Class Target

Capture Sol.

Decay Channel

Buncher

Phase Rotator

Cool-ing

Initial Cooling

Acceleration

0.2–1 GeV

1–5 GeV

Accelerators:
Single-Pass Linacs

μ Storage Ring

μ^+

ν

5 GeV

μ^-

ν

281m

Factory Goal:
10^{21} $^+$ & per year
within the accelerator acceptance

Share same complex

Muon Collider

Proton Driver

SC Linac

Accumulator

Buncher

Combiner

Front End

MW-Class Target

Capture Sol.

Decay Channel

Buncher

Phase Rotator

Cooling

Initial Cooling

Charge Separator

6D Cooling

μ^+

μ^-

Bunch Merge

6D Cooling

Final Cooling

Acceleration

Accelerators:
Linacs, RLA or FFAG, RCS

Collider Ring

E_{CoM}:
Higgs Factory
to
~10 TeV

μ^+ μ^-

-Collider Goals:
126 GeV ⇨
~14,000 Higgs/yr
Multi-TeV ⇨
Lumi > 10^{34}cm^{-2}s^{-1}

The Muon Cooling Section

studies for the arrangements of ion optical structure, absorber and rf section

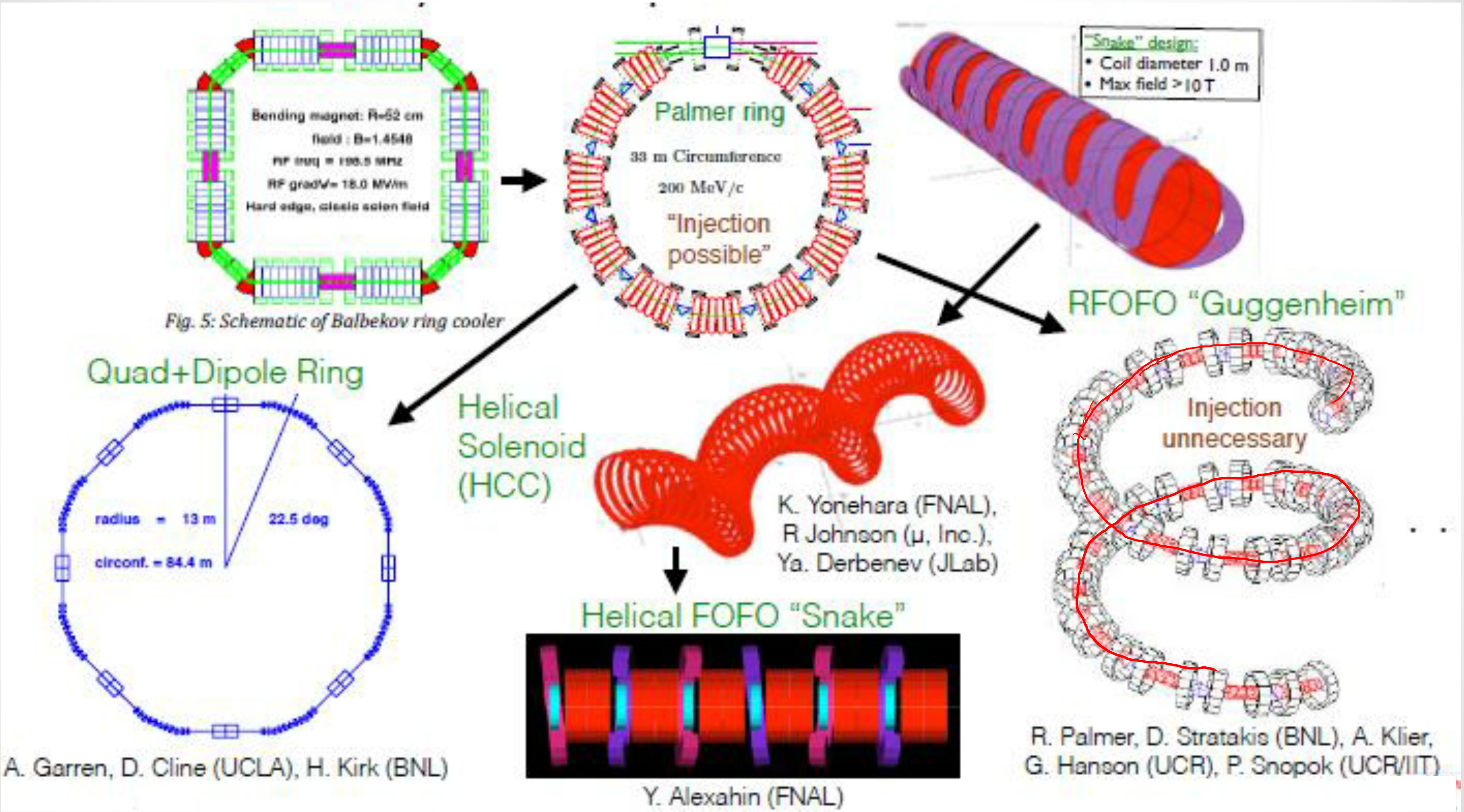


Fig. 5: Schematic of Balbekov ring cooler

A. Garren, D. Cline (UCLA), H. Kirk (BNL)

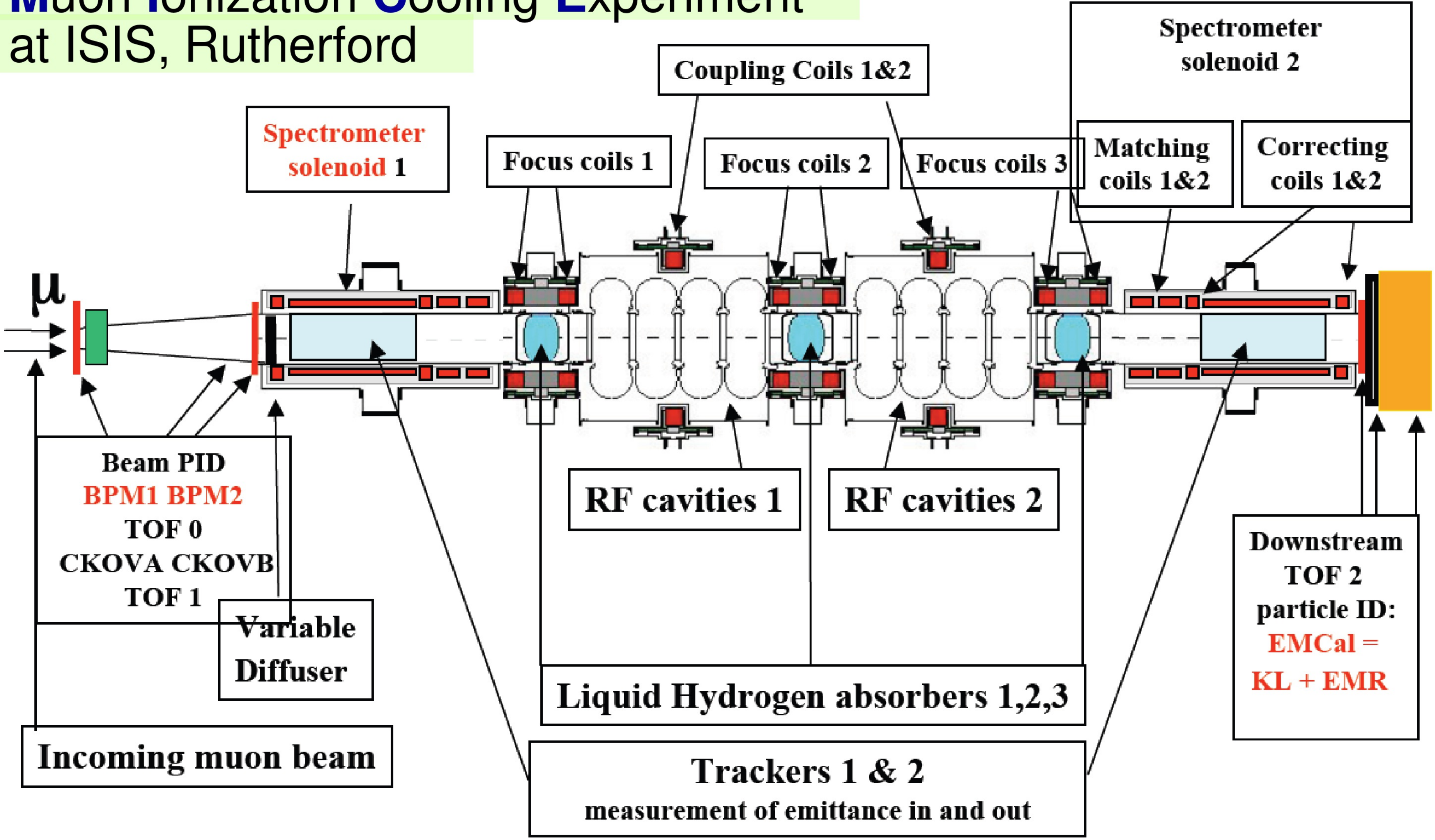
K. Yonehara (FNAL), R Johnson (μ, Inc.), Ya. Derbenev (JLab)

Y. Alexahin (FNAL)

R. Palmer, D. Stratakis (BNL), A. Klier, G. Hanson (UCR), P. Snopok (UCR/IIT)

MICE

Muon **I**onization **C**ooling **E**xperiment
at ISIS, Rutherford

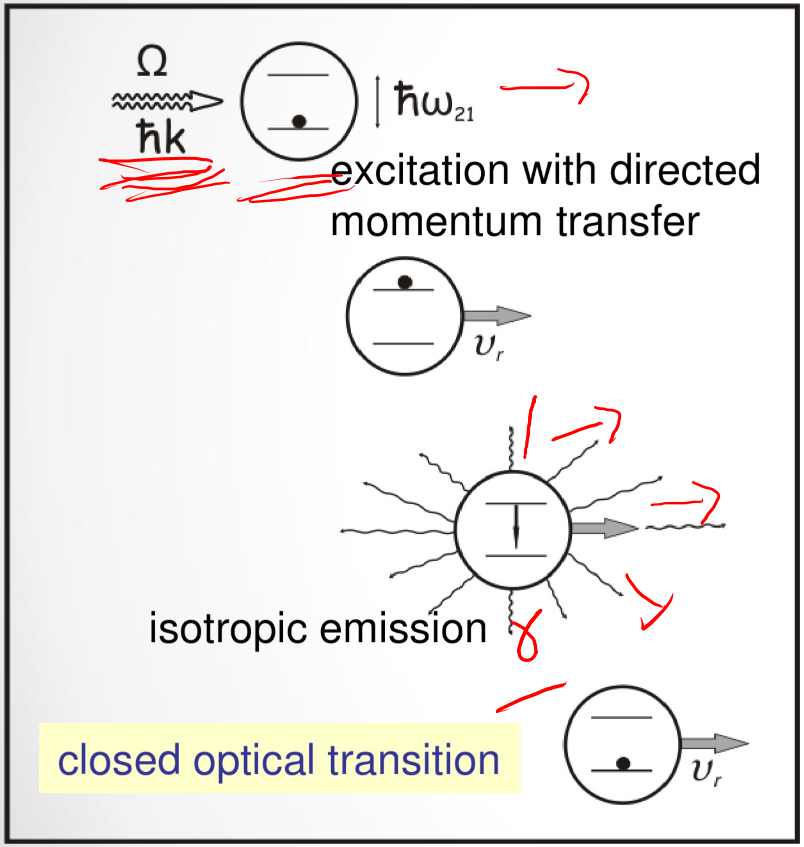


ion with e^-

3. Laser Cooling

↳ single-particle dynamics

$$\Omega = \gamma \omega_{21}(1-\beta\cos\theta)$$

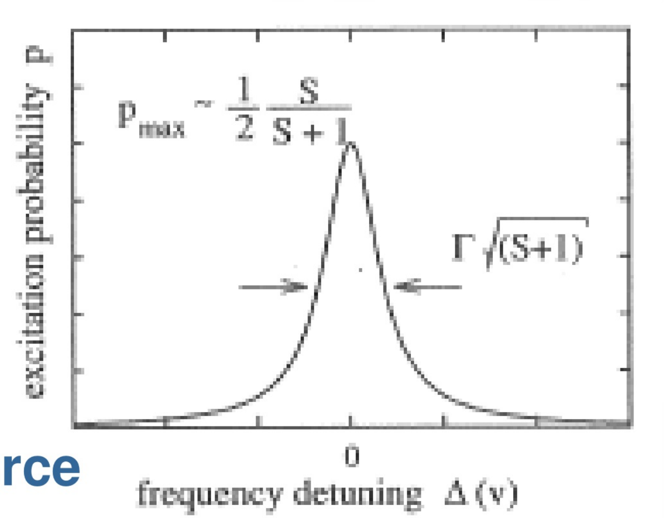


excitation with directed momentum transfer

isotropic emission

closed optical transition

the directed excitation and isotropic emission result in a transfer of velocity v_r



cooling force

$$\vec{F}(\vec{v},\vec{k}) = \frac{\hbar\vec{k}}{2} S\Gamma \frac{(\Gamma/2)^2}{(\omega-\omega_{21}-\vec{v}\vec{k}) + (\Gamma/2)^2(1+S)}$$

Lorentzian with width Γ/k ~ 10 m/s

minimum temperature $T_D = \frac{\hbar\Gamma}{2k_B}$ (Doppler limit)
typical 10^{-5} – 10^{-4} K

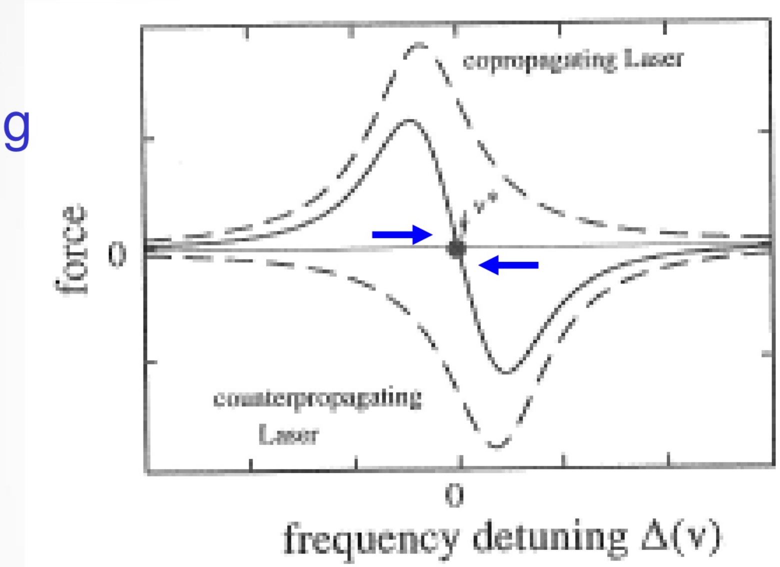
typical cooling time ~ 10 μs

drawback: only longitudinal cooling

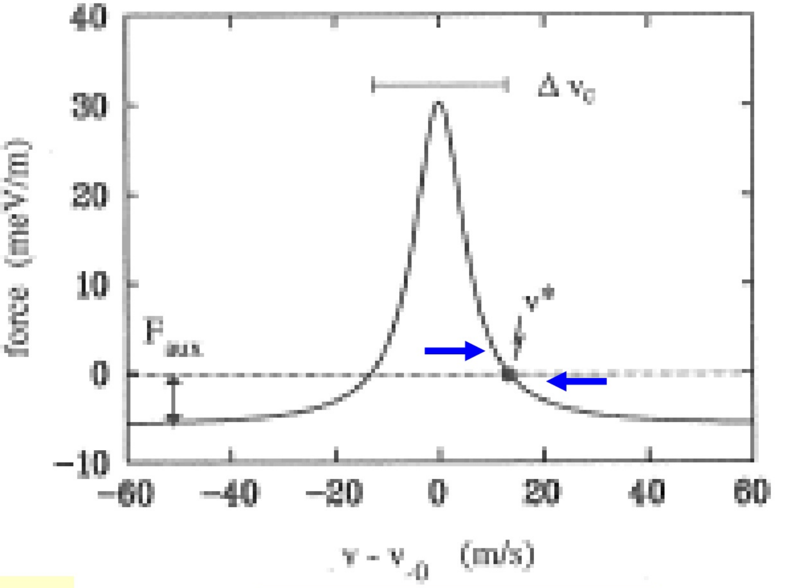
Laser Cooling

a single laser does not provide cooling (only acceleration or deceleration)

schemes
for cooling



two counter-propagating lasers (matched to beam velocity, but slightly detuned)



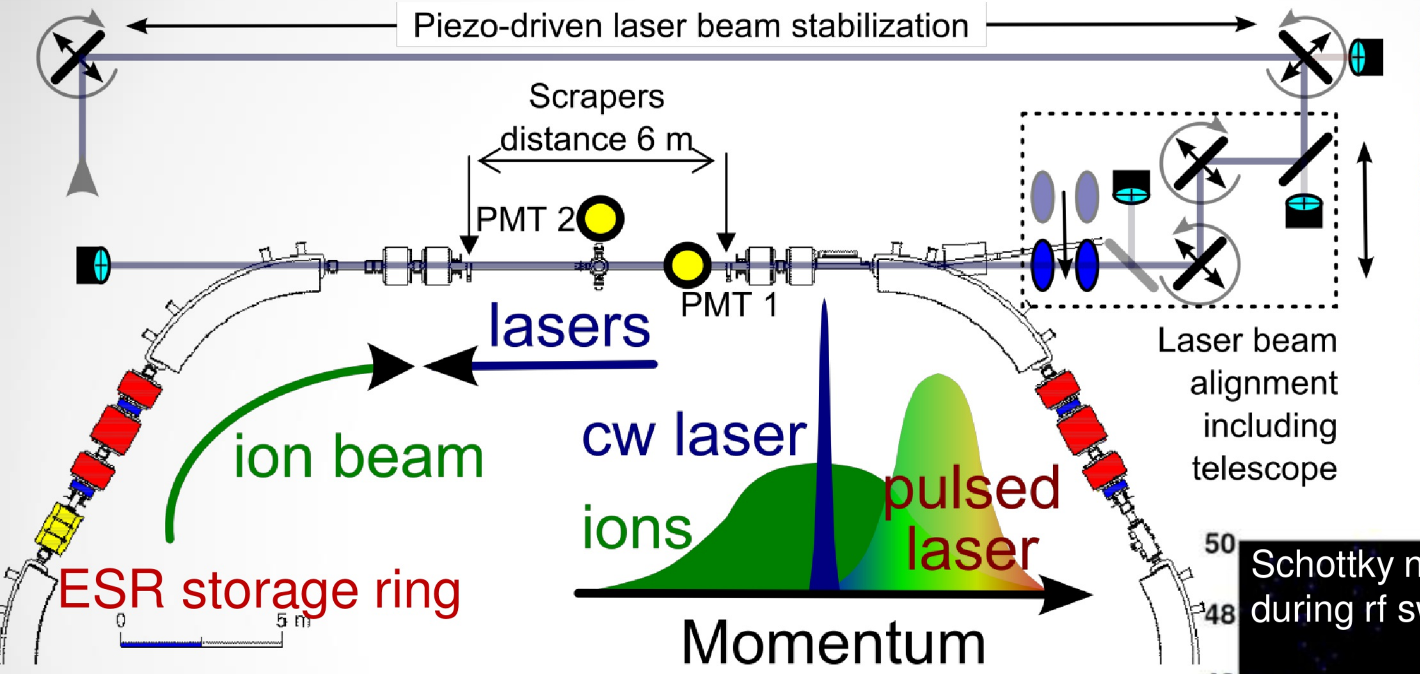
auxiliary force (betatron core, rf)

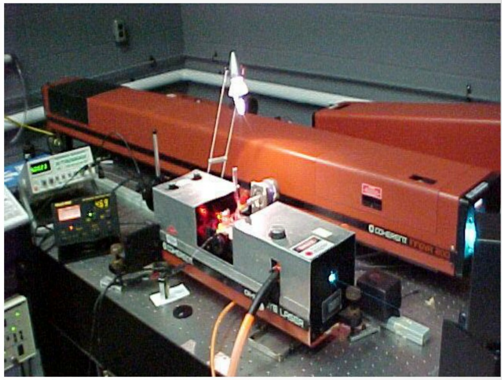
capture range of laser is limited ⇒ frequency sweep (snowplow)

ions studied so far: **$^7Li^{1+}$, $^9Be^{1+}$, $^{24}Mg^{1+}$, $^{12}C^{3+}$**

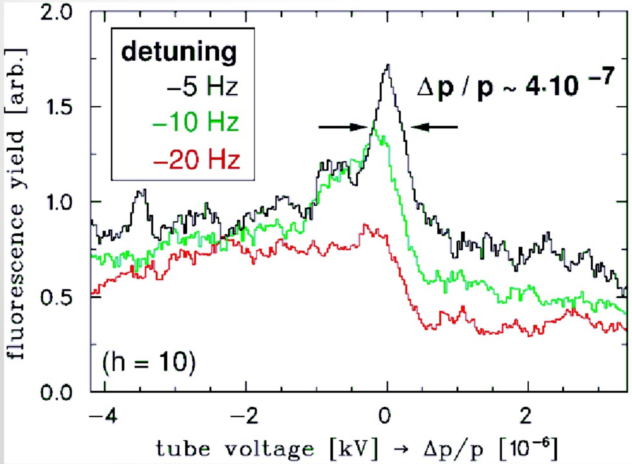
in future: Li-like heavy ions at relativistic energies
large relativistic energy ⇒ large excitation energy in PRF
Cooling rate increases with γ

Laser Cooling of C^{3+}

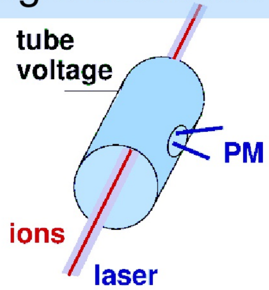




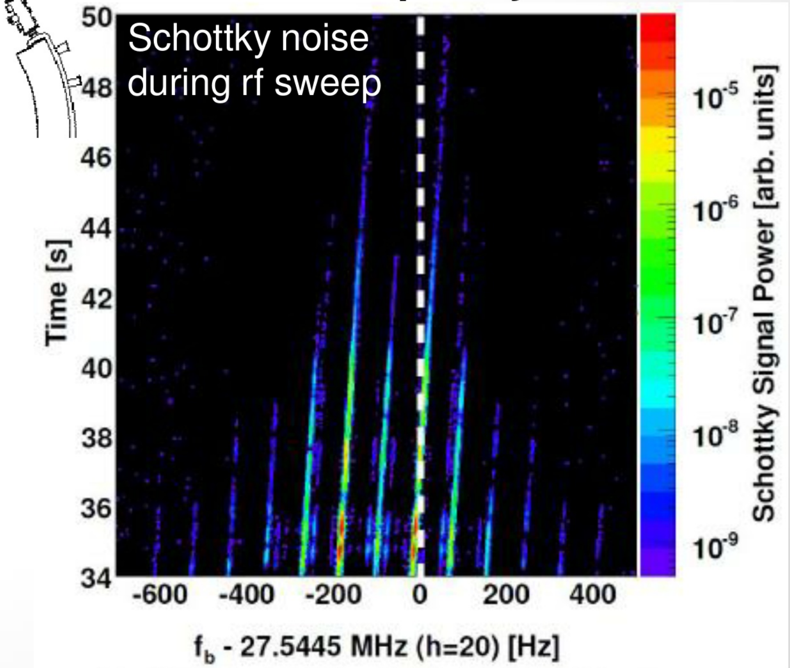
Argon ion laser (257.3 nm) frequency doubled



fluorescence light detection



probing the velocity distribution



4. Stochastic Cooling

⇒ Not single-particle dynamics

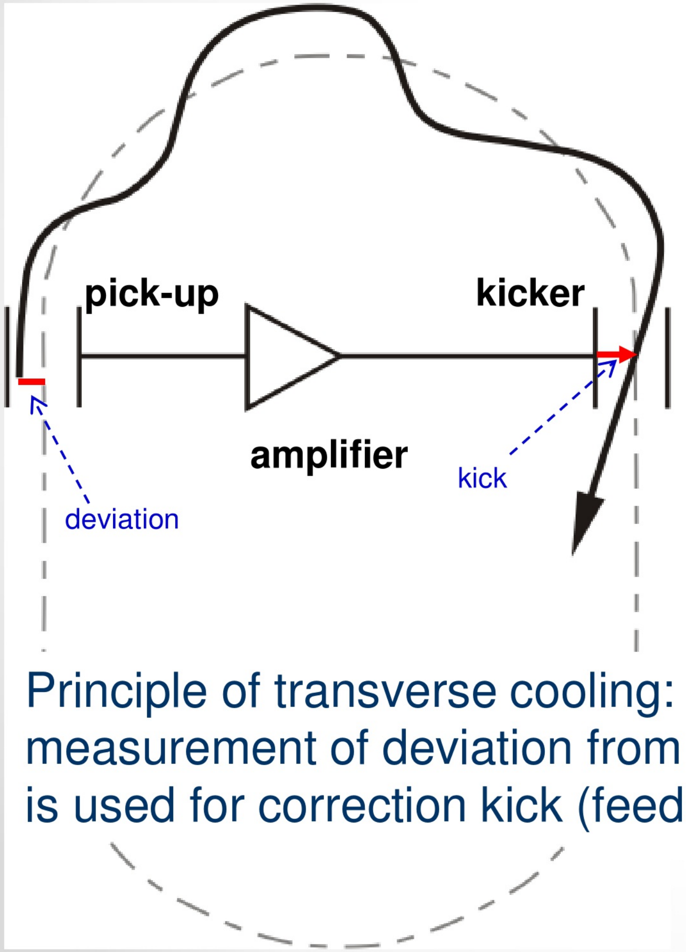
Finite N

$\langle x \rangle$ "Noisy"

$\propto 1/N^{1/2}$

Feedback

First cooling method which was successfully used for beam preparation



Principle of transverse cooling:
measurement of deviation from ideal orbit
is used for correction kick (feedback)

S. van der Meer, D. Möhl, L. Thorndahl et al.
(1925 – 2011) (1936-2012)

Conditions:

Betatron motion phase advance (pick-up to kicker): (n + ½) π

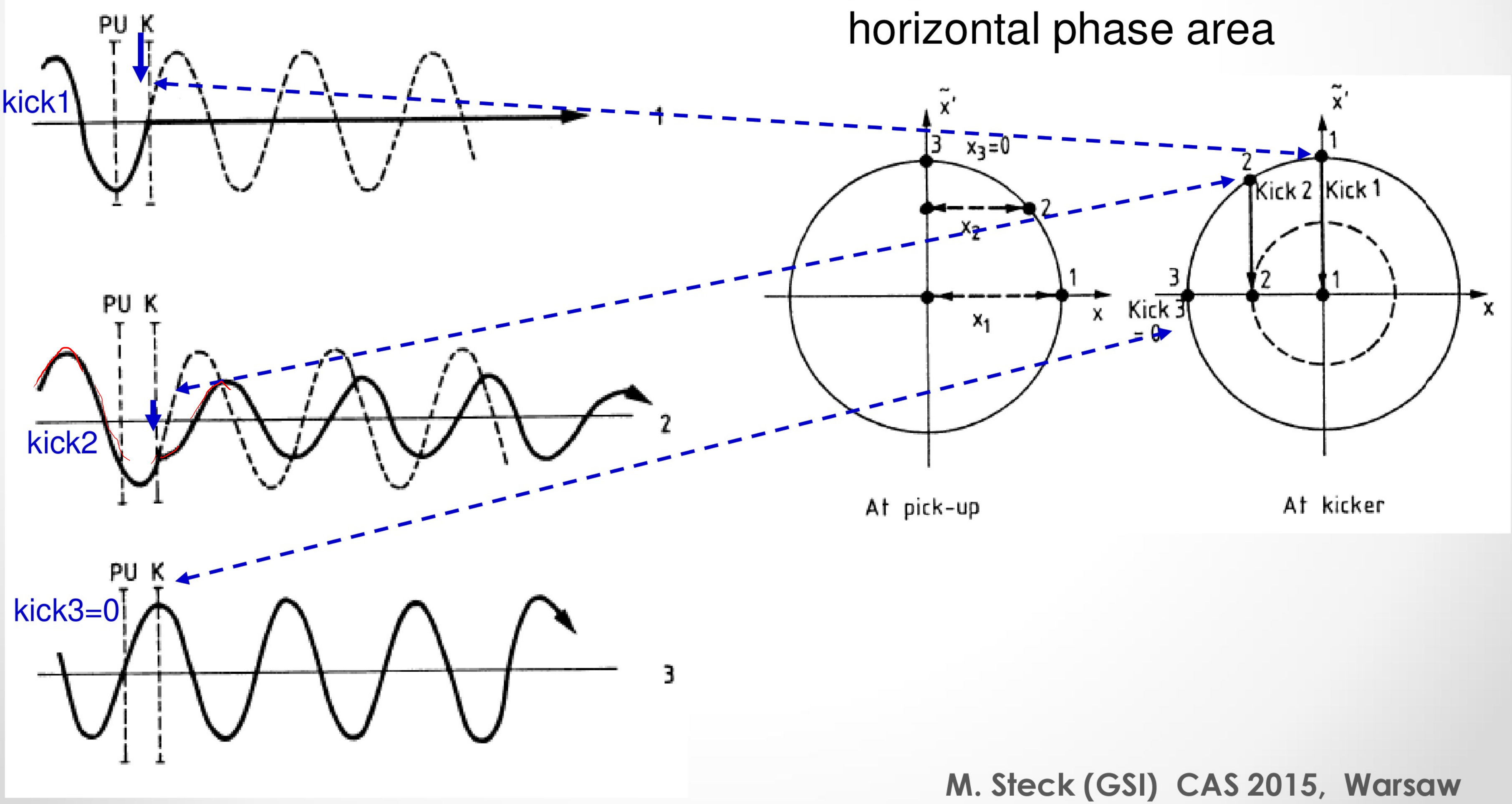
Signal travel time = time of flight of particle (between pick-up and kicker)

Sampling of sub-ensemble of total beam

Stochastic Cooling

single particle betatron motion along storage ring without (dashed) and with (full) correction kick

projection to two-dimensional horizontal phase area

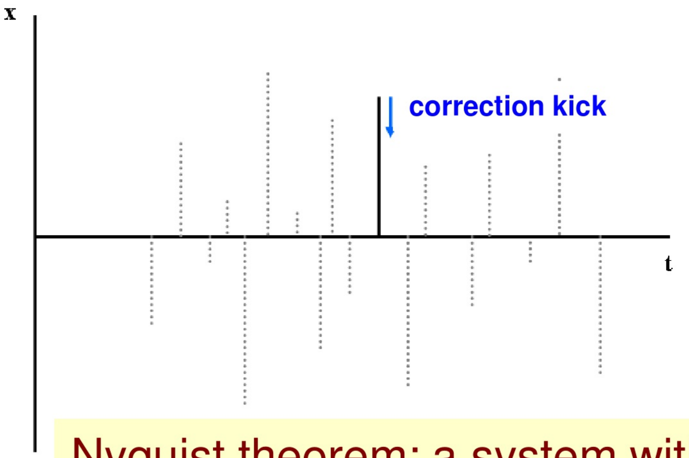


Stochastic Cooling

⇒ best tails or large amplitude

⇒ RHIC Au^{+79} ☺ $\sim 10^9$ particles/bunch

p^+ ☹ $\sim 10^{11}$ particles/bunch $100 \times \tau^{-1}$!

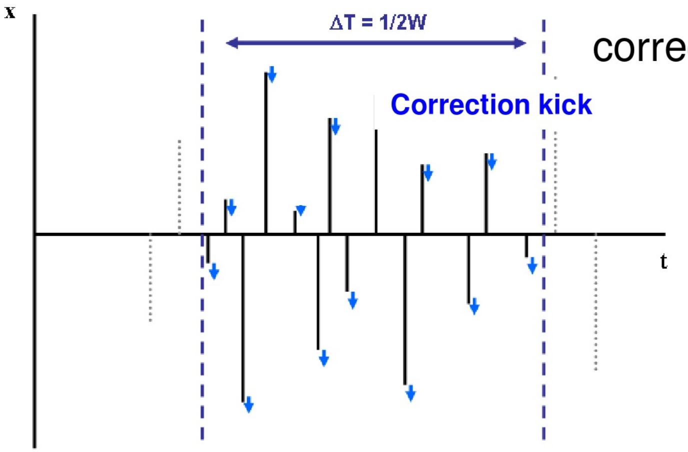


in time domain

correction kick
(unlimited resolution)

$$\Delta x = g \times x$$

Nyquist theorem: a system with a band-width Δf = W in frequency domain can resolve a minimum time duration ΔT=1/(2W)



correction kick $\Delta x = \frac{g}{N_s} \times \sum_{i=1..N_s} x_i \ , \quad N_s = N \frac{\Delta T}{T_0} = \frac{N}{2WT_0}$

For exponential damping ($x(t)=x(t_0)\cdot\exp(-(t-t_0)/\tau)$):

$$\tau^{-1} = T_0^{-1} \times \frac{\Delta x}{x} = \frac{g2W}{N}, \ if \sum_{i=1..N_s} x_i = x$$

cooling rate $\boxed{\tau^{-1} \leq \frac{2W}{N}}$ $if \ g \leq 1$

Stochastic Cooling

some refinements of cooling rate formula

noise: thermal or electronic noise adds to the beam signal

mixing: change of relative longitudinal position of particles due to momentum spread

cooling rate $$\lambda = \tau^{-1} = \frac{2W}{N}(\underbrace{2g}_{\text{cooling}} - \underbrace{g^2(M+U))}_{\text{heating}}$$

M mixing factor

U noise to signal ratio

maximum of cooling rate

$$\boxed{\lambda_{max} = \frac{2W}{N}\frac{1}{M+U}}$$

$$\frac{d\lambda}{dg} = 0 \Rightarrow g = \frac{1}{M+U}$$

further refinement (wanted ↔ unwanted mixing):

with wanted mixing M (kicker to pick-up)

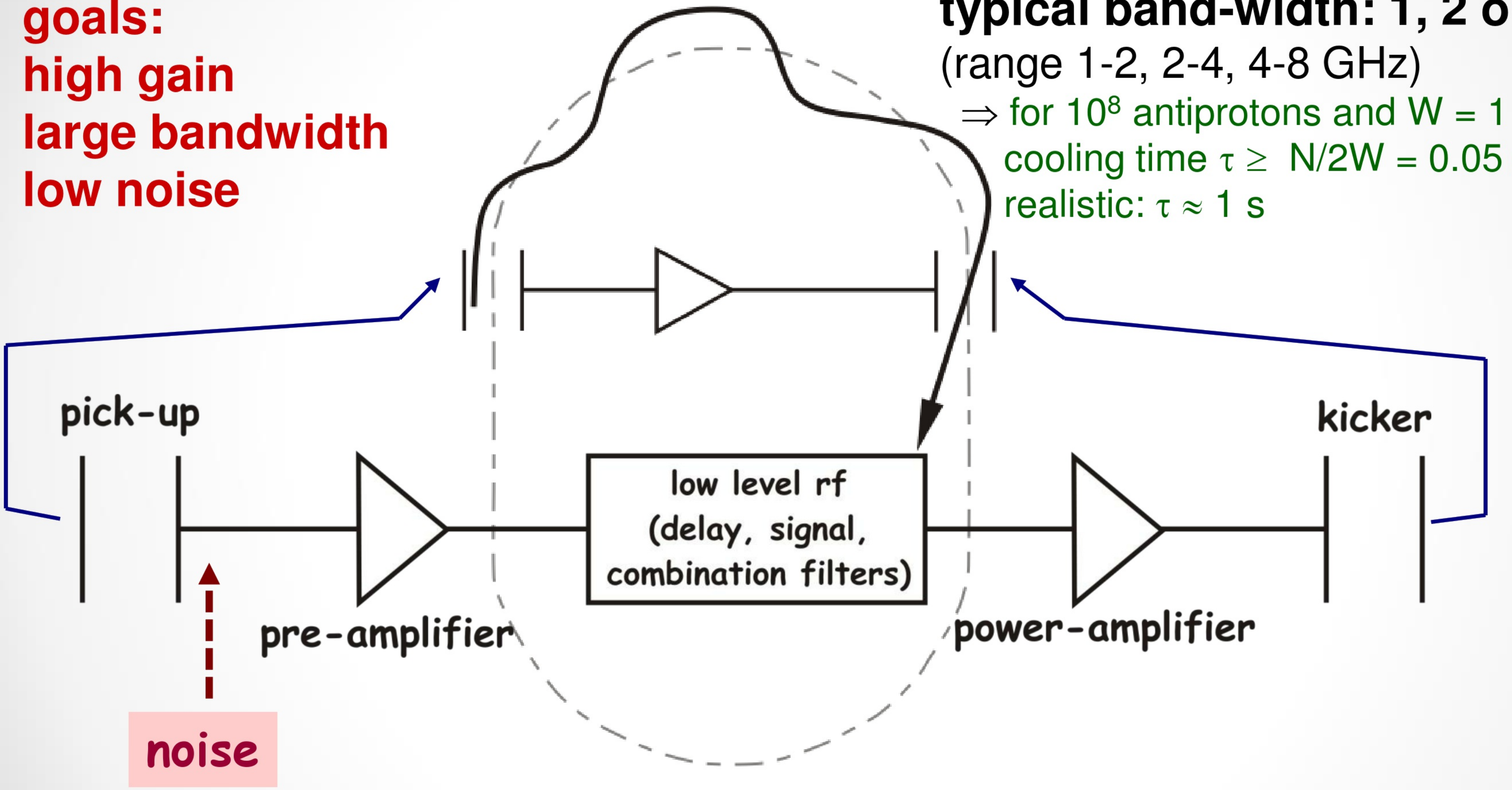
and unwanted mixing $\tilde{M}$ (pick-up to kicker)

$$\lambda = \tau^{-1} = \frac{2W}{N}(2g(1-\tilde{M}^2) - g^2(M+U))$$

Stochastic Cooling Circuit

goals:
high gain
large bandwidth
low noise

typical band-width: 1, 2 or 4 GHz
(range 1-2, 2-4, 4-8 GHz)
⇒ for 10^8 antiprotons and W = 1 GHz
cooling time $\tau \geq N/2W = 0.05$ s
realistic: $\tau \approx 1$ s



Transfer Function:

$$Z_{pick-up} \cdot G_{pick-up}(E) \cdot H(t_{delay}) \cdot F(E) \cdot G \cdot G_{kicker}(E) \cdot Z_{kicker}$$

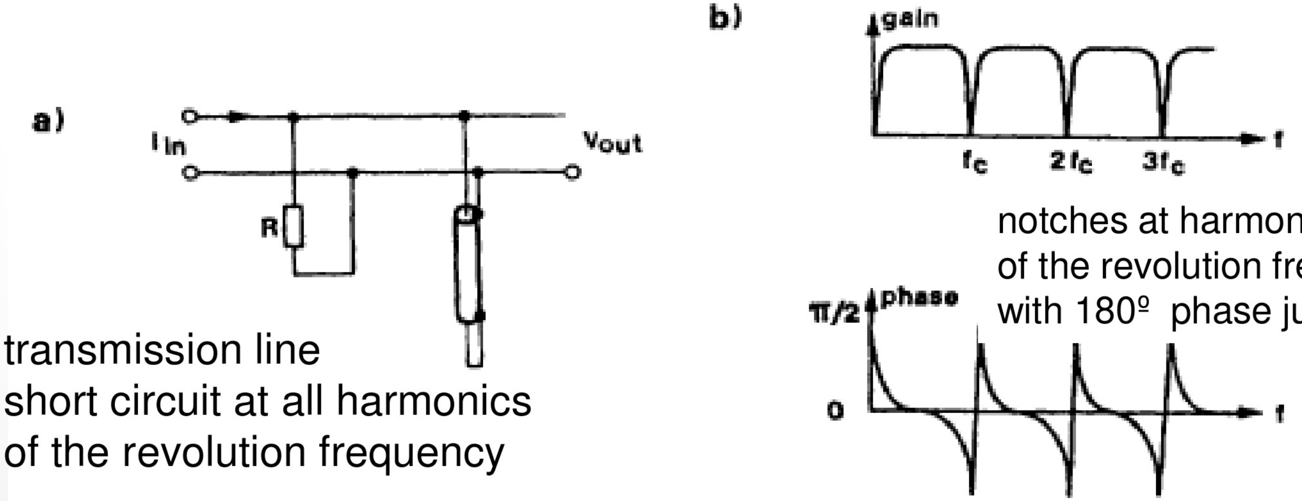
Longitudinal Stochastic Cooling

1) **Palmer cooling**

pick-up in dispersive section detects horizontal position
⇒ acceleration/deceleration kick corrects momentum deviation

2) **Notch filter cooling**

filter creates notches at the harmonics of the nominal revolution frequency
⇒ particles are forced to circulate at the nominal frequency

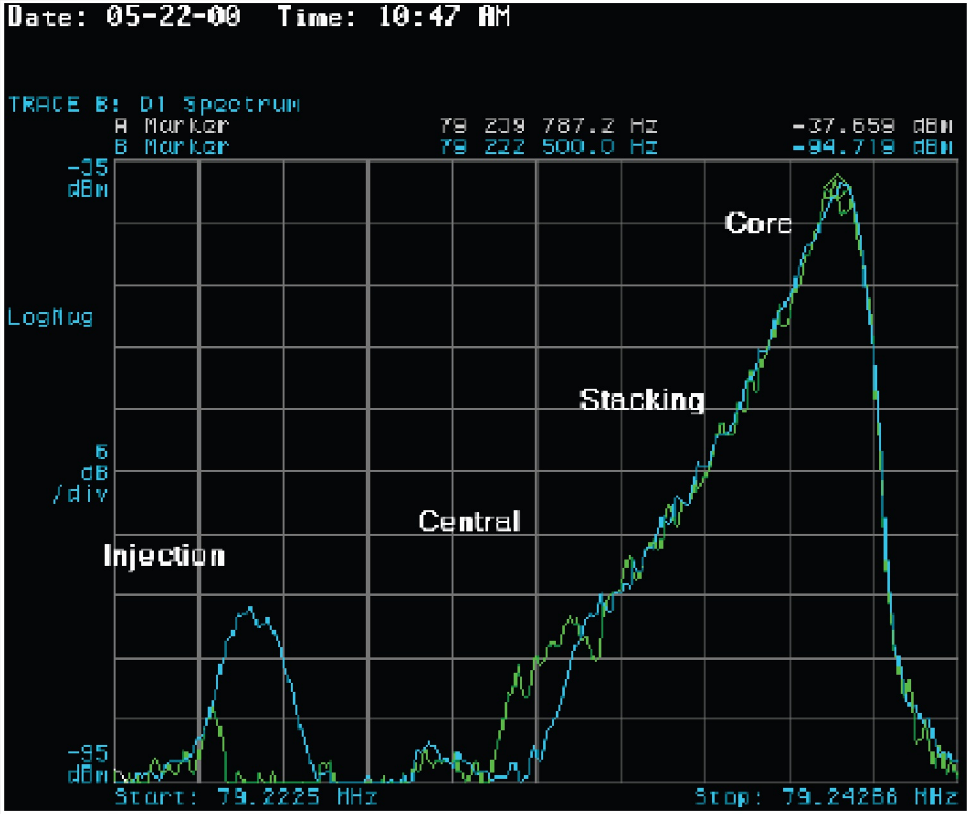


transmission line
short circuit at all harmonics
of the revolution frequency

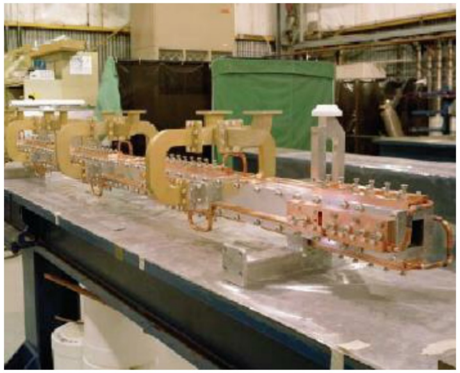
notches at harmonics
of the revolution frequency
with 180º phase jump

Antiproton Accumulation by Stochastic Cooling

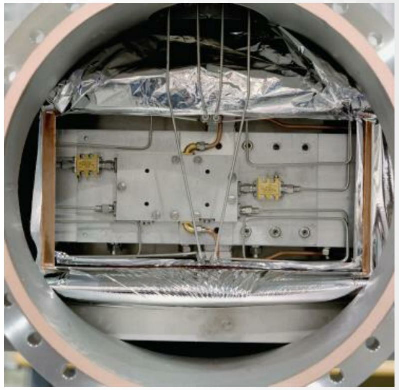
accumulation of 8 GeV antiprotons at accumulator ring, FNAL, shut down 09/2011
a similar facility AC/AA at CERN was shut down 11/1996



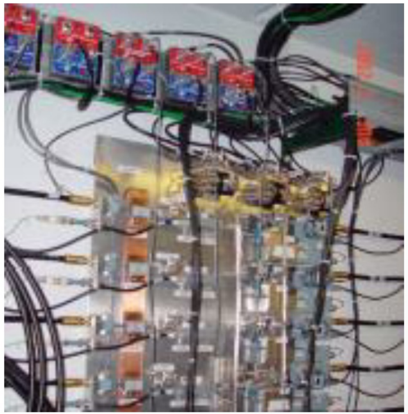
momentum distribution of accumulated antiproton beam



kicker array



cryogenic microwave amplifier



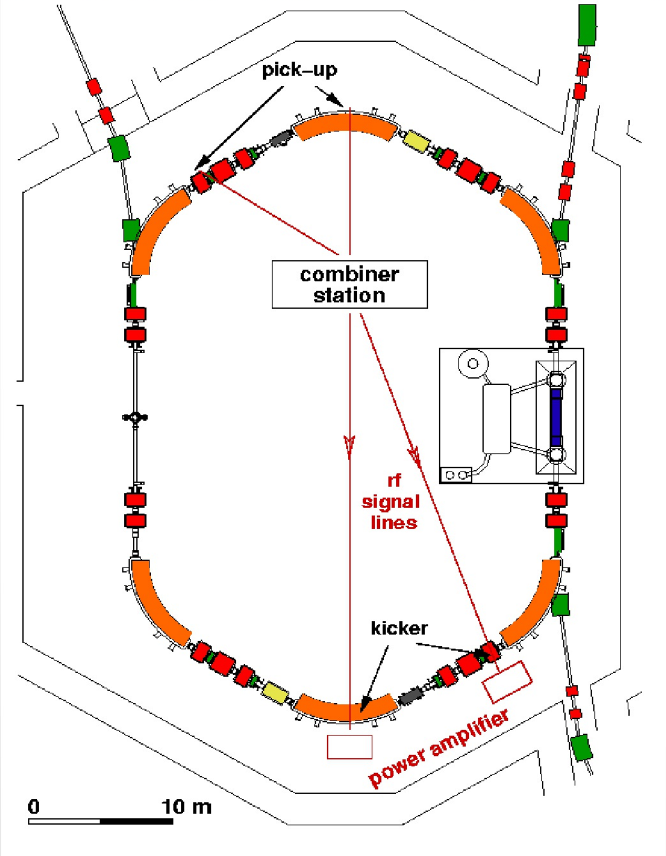
microwave electronics

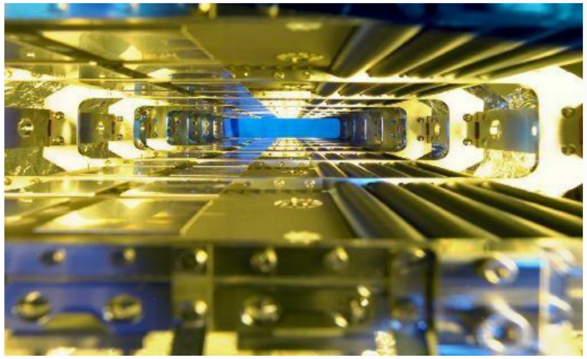
power amplifiers (TWTs)

Stochastic Cooling of Rare Isotopes at GSI

fast pre-cooling of hot fragment beams

energy 400 (-550) MeV/u
bandwidth 0.8 GHz (range 0.9-1.7 GHz)
$\delta p/p = \pm 0.35\ \%$ → $\delta p/p = \pm 0.01\ \%$
$\varepsilon = 10 \times 10^{-6}$ m → $\varepsilon = 2 \times 10^{-6}$ m





electrodes installed inside magnets

combination of signals from electrodes

power amplifiers for generation of correction kicks

Comparison of Cooling Methods

	Stochastic Cooling	Electron Cooling
Useful for:	low intensity beams	low energy
		all intensities
	hot (secondary) beams	warm beams (pre-cooled)
	high charge	high charge
	full 3D control	bunched beams
Limitations:	high intensity beams	space charge effects
/problems	beam quality limited	recombination losses
	bunched beams	high energy

laser cooling (of incompletely ionized ions) and ionization cooling (of muons) are quite particular and not general cooling methods

Trends in Beam Cooling

Stochastic cooling was mainly developed for the production of high intensity antiproton beams for colliders (CERN, FNAL, 1972 – 2011).
It is still in operation at AD (CERN), COSY (FZJ) and ESR (GSI).
It will also be used in the FAIR project (Germany) for cooling of antiprotons and rare isotope beams.

Demonstration of bunched beam stochastic cooling (2008) with heavy ions (BNL) and the achievement of increased luminosity made it very attractive for ion colliders.
Now it is proposed for the collider of the Russian NICA project.

Electron cooling was and still is used in low energy storage rings for protons, ions, secondary beams (antiprotons, rare isotopes).

Electron cooling is interesting for low energy storage rings, but also application at higher energies (MeV electron energies) is envisaged after the successful demonstration of the 4 MeV electron cooler at FNAL.
Bunched electron beam cooling and coherent electron cooling are in preparation for RHIC (BNL).

Muon (ionization) cooling is still far from implementation in a full scale machine.

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S. van der Meer, Rev. Mod. Phys. Vol. 57, No. 3 Part 1, 1985

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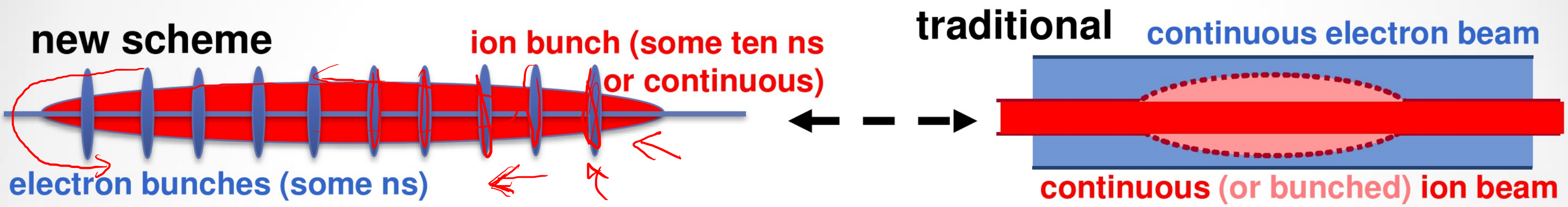
Ionization Cooling:

D. Neuffer, Introduction to Muon Cooling, Nucl. Instr. Meth. A 532 (2004) 26-31

Biannual Workshops on Beam Cooling: e. g. COOL’15, Jefferson Lab, USA

Bunched Beam Electron Cooling

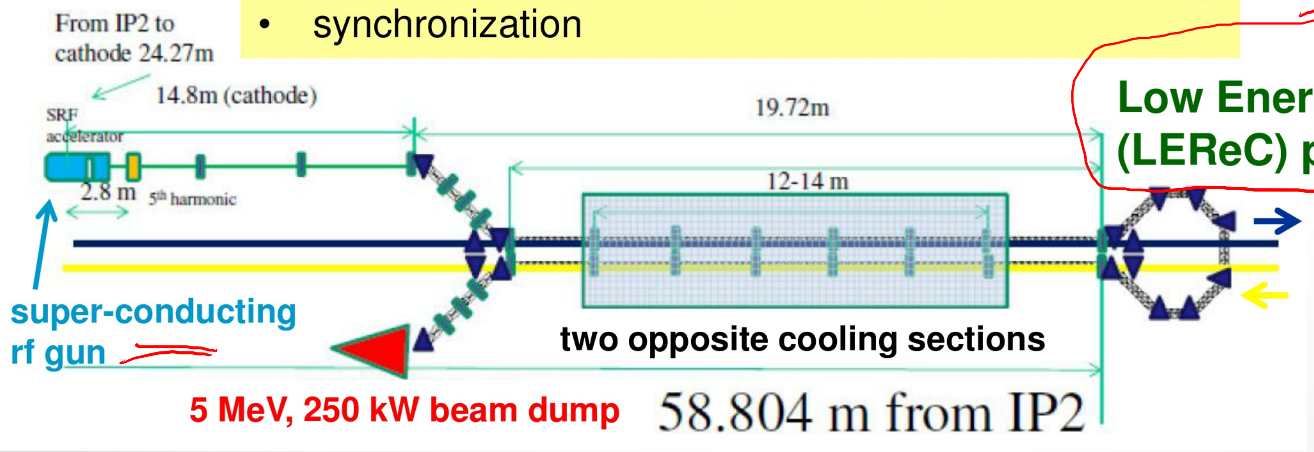
Electron cooling with electrostatic acceleration is limited in energy (5-10 MeV). A bunched electron beam offers the extension of the electron cooling method to higher energy (linear rf accelerator).



Issues:
- high intensity bunches (production, transport)
- momentum spread and emittance of bunches
- beam alignment
- magnetized ↔ non-magnetized (magnetic shielding)
- synchronization

M. Brucker

A. Fedotov & @ IMP



Low Energy RHIC e-Cooler (LEReC) project at BNL

Coherent Electron Cooling : EIC

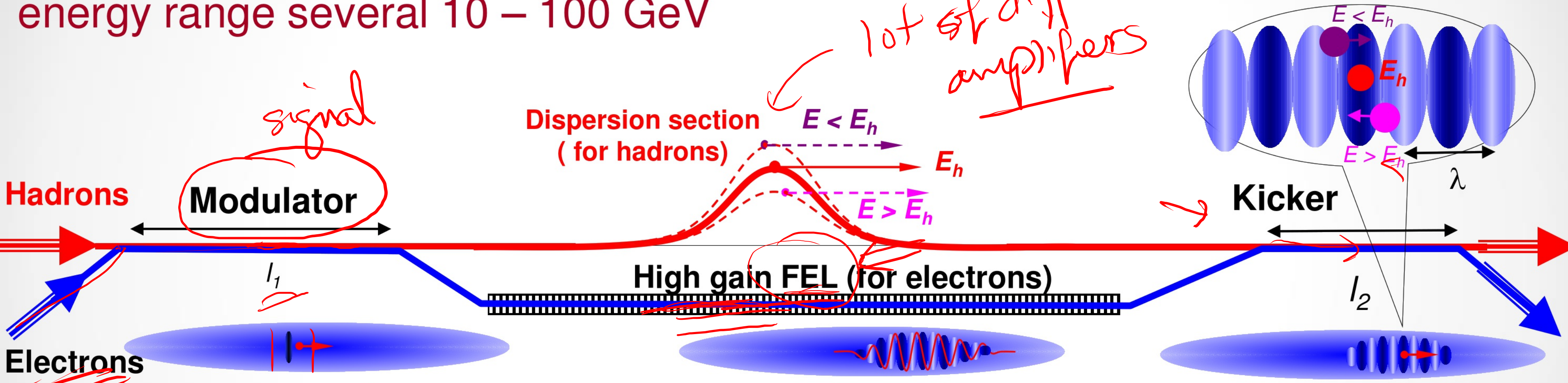
V. Litvinenko

A combination of electron and stochastic cooling concepts

proposed for fast cooling at highest energies

energy range several 10 – 100 GeV

lot of different amplifiers



quiet

Stochastic cooling with e⁻ beam as amplifier

- The Coherent Electron Cooling system has three major subsystems
  - **modulator**: the ions imprint a “density bump” on the electron distribution
  - **amplifier**: FEL interaction amplifies a density bump by orders of magnitude
  - **kicker**: the amplified & phase-shifted electron charge distribution is used to correct the velocity offset of the ions